

A Hybrid High-Order method for incremental associative plasticity with small deformations

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- Associative **plasticity** with **small** deformations
 - non-linear stress-strain constitutive relation (material nonlinearity)
 - history of the deformations (irreversible phenomena)
- Presence of **volumetric-locking** with primal H^1 -conforming formulation due to plastic incompressibility
- An alternative : using mixed methods but more unknowns, more expensive to build, saddle-point problem to solve ...

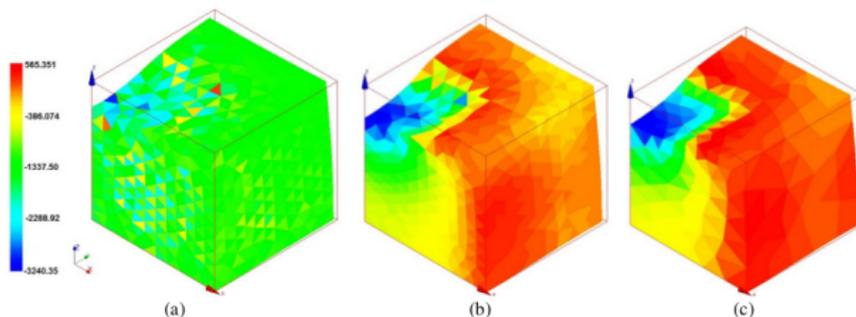


Figure 1 – Trace of the stress tensor for (a) P1 (b) P2 (c) P2/P1/P1

Main features of the proposed method

- **Primal** formulation
 - ⇒ More advantageous than mixed methods
- **Absence** of volumetric-locking
 - ⇒ More advantageous than primal methods
- Integration of the behavior law only at **cell-based** quadrature nodes
 - ⇒ More advantageous than discontinuous Galerkin (dG) methods
- Implementation in the open-source libraries *disk++* and *code_aster*
 - code : <https://github.com/wareHHOuse/diskpp>
 - code : <https://www.code-aster.org>

Some references for infinitesimal plasticity

Some references on **primal** formulations **without** volumetric-locking

- **discontinuous Galerkin (dG)**
 - [Noels, Radovitzsky 06]
 - [ten Eyck, Lew 06]
- **Hybrid Methods**
 - [Wulfinghoff, Bayat, Alipour, Reese 17]
 - [Krämer, Wieners, Wohlmuth, Wunderlich 16]
- **Virtual Element Method (VEM)**
 - [Chi, Beirão da Veiga, Paulino 17]
 - [Wriggers, Reddy, Rust, Hudobivnik 17]

Key ideas of Hybrid High-Order (HHO) methods

- Primal formulation with **cells** and **faces** unknowns
- **Local reconstruction and stabilization**
 - Symmetric gradient tensor field reconstructed in $\mathbb{P}_d^k(T; \mathbb{R}_{\text{sym}}^{d \times d})$
 - Stabilization connecting cell and faces unknowns
- References
 - diffusion problem [Di Pietro, Ern, Lemaire, CMAM 14]
 - quasi-incompressible linear elasticity [Di Pietro, Ern, CMAME 15]
 - nonlinear elasticity with small def. [Botti, Di Pietro, Sochala, SINUM 17]
 - hyperelasticity with finite deformations [Abbas, Ern, Pignet, CM 18]

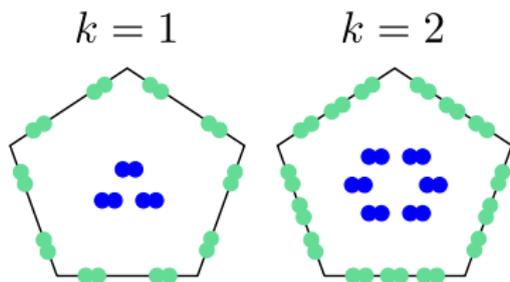


Figure 2 – Face (green) and Cell (blue) unknowns

- Support of **polytopal meshes** (with possibly nonconforming interfaces)
- **Arbitrary approximation order** $k \geq 1$
 - h^{k+1} convergence in energy-norm
 - h^{k+2} convergence in L^2 -norm with elliptic regularity
- Dimension-**independent** construction
- **Attractive** computational costs
 - Compact stencil (only neighbourhood faces)
 - Cell unknowns are eliminated locally by static condensation
 - Reduced size $N_{dofs}^{hho} \approx k^2 \text{card}(\mathcal{F}^h)$ vs. $N_{dofs}^{dG} \approx k^3 \text{card}(\mathcal{T}^h)$
- Local principle of virtual work (**equilibrated tractions**)
- HHO methods are **bridged** to HDG and ncVEM
 - [Cockburn, Di Pietro, Ern 16]

- Let $\Omega_0 \in \mathbb{R}^d$ ($d=2,3$), be a bounded connected polytopal domain
- Let $\underline{\mathbf{f}}$ and $\underline{\mathbf{t}}$ be given volumetric and surface (on Γ_n) loads
- Let $\underline{\mathbf{u}}_D$ be a given imposed displacement on Γ_d
- **History** of the deformations : \rightarrow we introduce the internal state variables $\underline{\chi}$
- For all $1 \leq n \leq N$, find $\underline{\mathbf{u}}^n \in V_D := \{\underline{\mathbf{v}} \in H^1(\Omega_0; \mathbb{R}^d) \mid \underline{\mathbf{v}} = \underline{\mathbf{u}}_D \text{ on } \Gamma_d\}$ s.t.

$$\int_{\Omega_0} \underline{\underline{\boldsymbol{\sigma}}}(\underline{\mathbf{u}}^n) : \underline{\underline{\boldsymbol{\varepsilon}}}(\underline{\mathbf{v}}) d\Omega_0 = \int_{\Omega_0} \underline{\mathbf{f}}^n \cdot \underline{\mathbf{v}} d\Omega_0 + \int_{\Gamma_n} \underline{\mathbf{t}}^n \cdot \underline{\mathbf{v}} d\Gamma \text{ for all } \underline{\mathbf{v}} \in V_0,$$

and

$$\underline{\underline{\boldsymbol{\sigma}}}(\underline{\mathbf{u}}^n) = \text{PLASTICITY}(\underline{\chi}, \underline{\underline{\boldsymbol{\varepsilon}}}(\underline{\mathbf{u}}^{n-1}), \underline{\underline{\boldsymbol{\varepsilon}}}(\underline{\mathbf{u}}^n)).$$

where PLASTICITY is a generic behavior integrator

Local DOFs space

- Let $M^h := (\mathcal{T}^h, \mathcal{F}^h)$ be a mesh of Ω_0 with \mathcal{T}^h the set of cells and \mathcal{F}^h the set of (planar) faces
- Let a polynomial degree $k \geq 1$, for all $T \in \mathcal{T}^h$

$$(\underline{\mathbf{v}}_T, \underline{\mathbf{v}}_{\partial T}) \in \underbrace{\mathbb{P}_d^k(T; \mathbb{R}^d)}_{\text{local cell dofs}} \times \underbrace{\mathbb{P}_{d-1}^k(\mathcal{F}_{\partial T}; \mathbb{R}^d)}_{\text{local faces dofs}}.$$

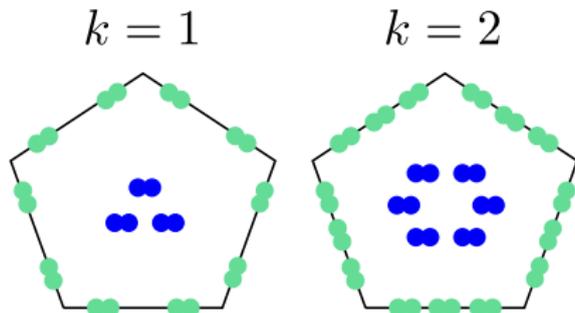


Figure 3 – Local DOFs for $k = 1, 2$. Cell unknowns are eliminated by static condensation

$$\underline{\underline{\mathbf{E}}}_T^k : \mathbb{P}_d^k(T; \mathbb{R}^d) \times \mathbb{P}_{d-1}^k(\mathcal{F}_{\partial T}; \mathbb{R}^d) \rightarrow \underbrace{\mathbb{P}_d^k(T; \mathbb{R}_{\text{sym}}^{d \times d})}_{\text{local strain space}}$$

- The reconstructed strain $\underline{\underline{\mathbf{E}}}_T^k(\underline{\mathbf{v}}_T, \underline{\mathbf{v}}_{\partial T})$ solves, $\forall \underline{\underline{\boldsymbol{\tau}}} \in \mathbb{P}_d^k(T; \mathbb{R}_{\text{sym}}^{d \times d})$

$$(\underline{\underline{\mathbf{E}}}_T^k(\underline{\mathbf{v}}_T, \underline{\mathbf{v}}_{\partial T}), \underline{\underline{\boldsymbol{\tau}}})_{\underline{\underline{\mathbf{L}}}(T)} = (\underline{\underline{\nabla}}^{\text{sym}} \underline{\mathbf{v}}_T, \underline{\underline{\boldsymbol{\tau}}})_{\underline{\underline{\mathbf{L}}}(T)} + (\underline{\mathbf{v}}_{\partial T} - \underline{\mathbf{v}}_T, \underline{\underline{\boldsymbol{\tau}}} \underline{\mathbf{n}}_T)_{\underline{\underline{\mathbf{L}}}(\partial T)}.$$

- well-posed local Neumann problem
- local **scalar** mass-matrix of size $\binom{k+d}{k}$ to invert (ex : $k = 2, d = 3, \text{size} = 10$)
- $\underline{\underline{\mathbf{E}}}_T^k$ depends only on the **geometry** of T (for k fixed)

Stabilization operator

- **Problem** : $\underline{\underline{\mathbf{E}}}_T^k(\underline{\mathbf{v}}_T, \underline{\mathbf{v}}_{\partial T}) = \underline{\underline{\mathbf{0}}} \not\Rightarrow \underline{\mathbf{v}}_T = \underline{\mathbf{v}}_{\partial T} = \text{cste}$
⇒ We have to "connect" the cell and faces unknowns

- Hence, we penalize the **difference between the faces unknowns and the trace of the cell unknowns** : $\underline{\mathbf{S}}_{\partial T}^k(\underline{\mathbf{v}}_{\partial T} - \underline{\mathbf{v}}_{T|\partial T}) \in \mathbb{P}_{d-1}^k(\mathcal{F}_{\partial T}; \mathbb{R}^d)$,

$$\underline{\mathbf{S}}_{\partial T}^k(\underline{\mathbf{v}}_{\partial T} - \underline{\mathbf{v}}_{T|\partial T}) = \underbrace{\underline{\mathbf{\Pi}}_{\partial T}^k(\underline{\mathbf{v}}_{\partial T} - \underline{\mathbf{v}}_{T|\partial T})}_{\text{HDG term}} - \underbrace{(\underline{\mathbf{I}}_d - \underline{\mathbf{\Pi}}_T^k)\underline{\mathbf{D}}_T^{k+1}(\underline{\mathbf{0}}, \underline{\mathbf{v}}_{\partial T} - \underline{\mathbf{v}}_{T|\partial T})}_{\text{high-order correction}}$$

where $\underline{\mathbf{\Pi}}_{\partial T}^k$ is the L^2 -projector on ∂T , $\underline{\mathbf{\Pi}}_T^k$ the L^2 -projector on T , and $\underline{\mathbf{D}}_T^{k+1}$ is a reconstructed displacement field

- Different to the HDG-stabilization operator
- The **high-order** correction is a distinctive feature of HHO methods

Global discrete problem

For all $1 \leq n \leq N$, find $(\underline{\mathbf{u}}_{\mathcal{T}^h}^n, \underline{\mathbf{u}}_{\mathcal{F}^h}^n) \in \left\{ \prod_{T \in \mathcal{T}^h} \mathbb{P}_d^k(T; \mathbb{R}^d) \right\} \times \left\{ \prod_{F \in \mathcal{F}^h} \mathbb{P}_{d-1}^k(F; \mathbb{R}^d) \right\}$

$$\begin{aligned} & \sum_{T \in \mathcal{T}^h} (\underline{\boldsymbol{\sigma}}(\underline{\mathbf{u}}_T^n, \underline{\mathbf{u}}_{\partial T}^n), \underline{\mathbf{E}}_T^k(\delta \underline{\mathbf{v}}_T, \delta \underline{\mathbf{v}}_{\partial T}))_{\underline{\mathbf{L}}^2(T)} \\ & + \sum_{T \in \mathcal{T}^h} \beta h_T^{-1} (\underline{\mathbf{S}}_{\partial T}^k(\underline{\mathbf{u}}_{\partial T}^n - \underline{\mathbf{u}}_{T|\partial T}^n), \underline{\mathbf{S}}_{\partial T}^k(\delta \underline{\mathbf{v}}_{\partial T} - \delta \underline{\mathbf{v}}_{T|\partial T}))_{\underline{\mathbf{L}}^2(\partial T)} \\ & = \sum_{T \in \mathcal{T}^h} (\underline{\mathbf{f}}, \delta \underline{\mathbf{v}}_T)_{\underline{\mathbf{L}}^2(T)} + \sum_{F \in \mathcal{F}_{b,n}^h} (\underline{\mathbf{t}}, \delta \underline{\mathbf{v}}_F)_{\underline{\mathbf{L}}^2(F)}, \quad \forall (\delta \underline{\mathbf{v}}_{\mathcal{T}^h}, \delta \underline{\mathbf{v}}_{\mathcal{F}^h}). \end{aligned}$$

and for all quadrature points

$$\underline{\boldsymbol{\sigma}}(\underline{\mathbf{u}}_T^n, \underline{\mathbf{u}}_{\partial T}^n) = \text{PLASTICITY}(\underline{\boldsymbol{\chi}}_T, \underline{\mathbf{E}}_T^k(\underline{\mathbf{u}}_T^{n-1}, \underline{\mathbf{u}}_{\partial T}^{n-1}), \underline{\mathbf{E}}_T^k(\underline{\mathbf{u}}_T^n, \underline{\mathbf{u}}_{\partial T}^n))$$

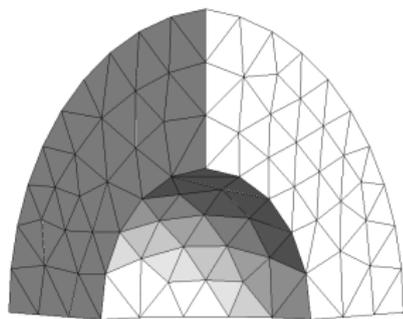
with $\beta \simeq 2\mu$ an user-dependent stabilization parameter

Numerical examples

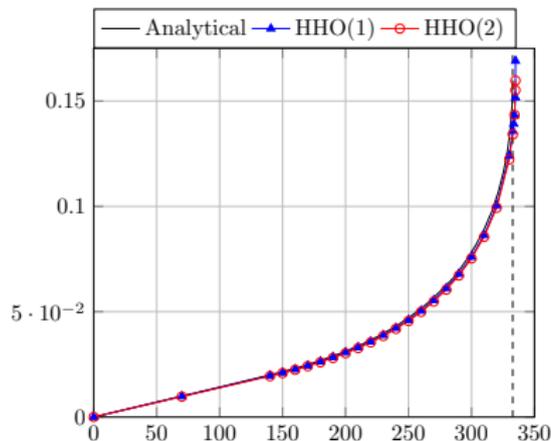
- **Nonlinear** problem to solve (material nonlinearity)
- Iterative resolution with a **Newton method**
- Static condensation performed at **each Newton's iteration**
- **Offline** computations (gradient and stabilization operators precomputed)
- Implementation in the open-source libraries *disk++* and *code_aster*
- Verification on analytical solution :
 - **Absence of volumetric-locking** due to plastic incompressibility
- Comparison to P^2 and $P^2/P^1/P^1$ (UPG) methods

Numerical example : sphere under internal pressure I

- Perfect J_2 -plasticity
- Increase the internal pressure until the limit load

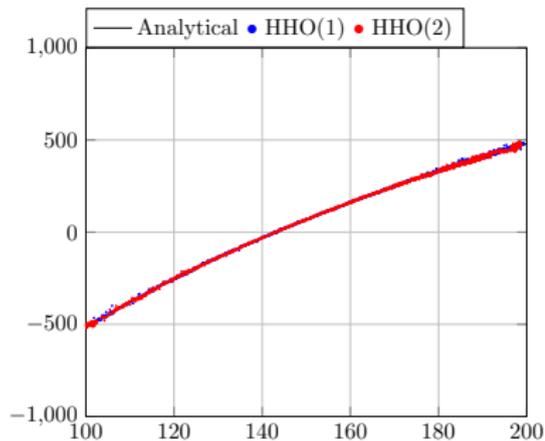


(a) Mesh.

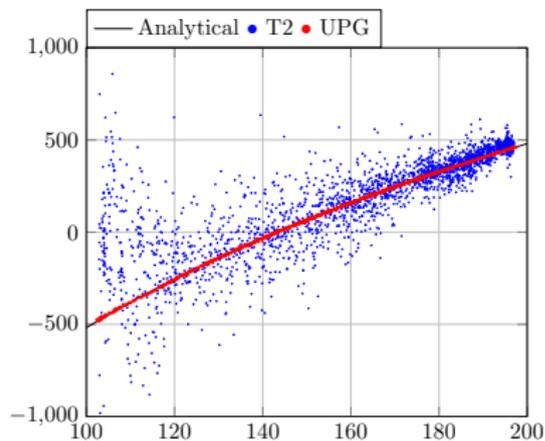


(b) Radial displacement vs. internal pressure.

Numerical example : sphere under internal pressure II



(c)



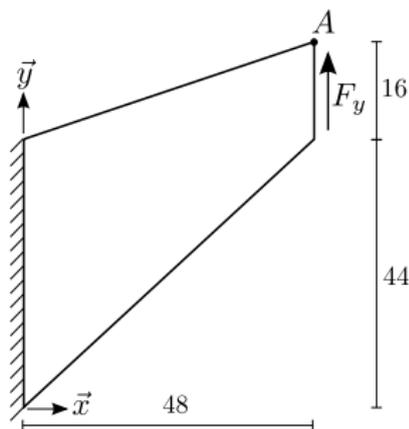
(d)

Figure 4 – Trace of the stress tensor at the quadrature points at the limit load

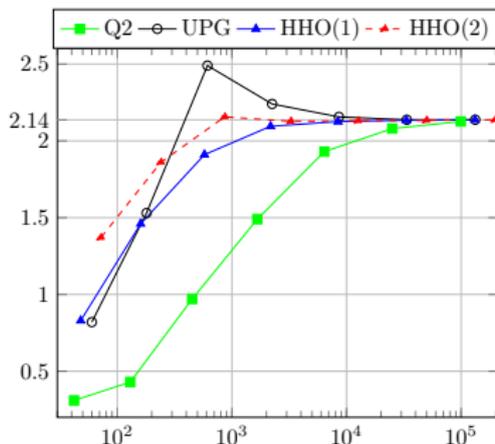
- Absence of volumetric-locking for HHO and mixed (UPG) methods

Numerical example : quasi-incompressible Cook's membrane

- Linear isotropic hardening with J_2 -plasticity ($\nu = 0.4999$)



(a) Geometry and BC.

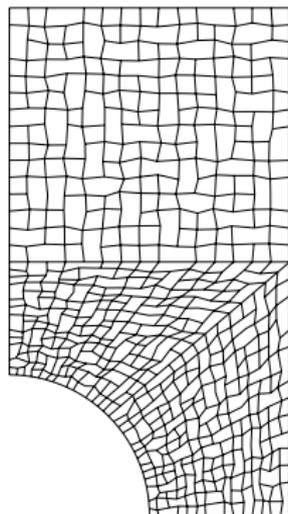


(b) Displacement of A vs. #dofs.

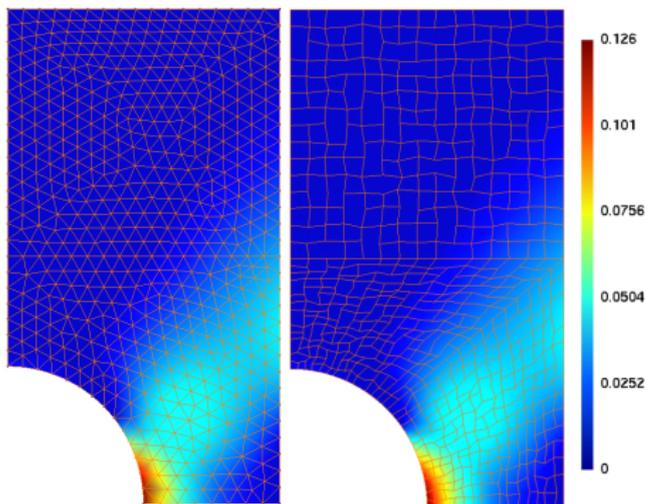
- HHO(2) **out-performs** compare to mixed (UPG) methods (same displacement order)

Perforated strip under uniaxial extension

- Combined linear kinematic and isotropic hardening with J_2 -plasticity



(c) Polygonal mesh



(d) Equivalent plastic strain with HHO(2)

- Support of **polyhedral** meshes

Unstabilized HHO method on simplicial meshes

- **How** to choose the stabilization parameter β :
 - $\beta = 2\mu$ is the value for the **linear elasticity**
 - **No general** theory on the choice of β
 - If β is too small \Rightarrow **Difficulties** to converge
 - If β is too large \Rightarrow The system is **ill-conditioned**
- Original idea for dG : [John, Neilan, Smears 16]
 - Based on the properties of the **Raviart–Thomas space**
- Gradient reconstruction in $\mathbb{P}_d^{k+1}(T; \mathbb{R}_{sym}^{d \times d})$ (larger space)
 - ex : $k = 2, d = 3$, size = 20 for $\mathbb{P}_d^{k+1}(T; \mathbb{R}^{d \times d})$ vs 10 for $\mathbb{P}_d^k(T; \mathbb{V}^{d \times d})$
- **No additional** stabilization is needed ($\beta \equiv 0$)
- **Lower** convergence rates (h^k in energy-norm and h^{k+1} in L^2 -norm)
- **Comparable** numerical cost vs. stabilized HHO methods

Conclusions and perspectives

- Conclusion :
 - Adaptation of HHO methods to associative **plasticity** with small deformations
 - **Primal** formulation
 - **Absence** of volumetric-locking
- Perspectives of this work :
 - Extension to finite plasticity
 - Introduction of contact and friction using Nitsche's method (with F. Chouly)

Thank you for your attention

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Reference : M. Abbas, A. Ern and NP, "A Hybrid High-Order method for incremental associative plasticity with small deformations", Comput. Methods. Appl. Mech. Engrg. (2018)