

# Microscopic Hamiltonian dynamics perturbed by a conservative noise

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# Introduction

- Fourier's law : Consider a macroscopic system in contact with two heat baths with different temperatures  $T_\ell \neq T_r$ . When the system reaches its steady state  $\langle \cdot \rangle_{ss}$ , one expects Fourier's law holds:

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- $J(q)$  is the energy current;  $T(q)$  the local temperature;  $\kappa(T)$  the conductivity.
- If system has (microscopic) size  $N$ , finite conductivity means  $\langle J \rangle_{ss} \sim N^{-1}$ .

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- Non linearity is extremely important to have normal heat conduction.
- But it is not sufficient : It has been observed experimentally and numerically for nonlinear chains that if  $d \leq 2$  and momentum is conserved ( $\Leftrightarrow W = 0$ , unpinned) then conductivity is still infinite (finite otherwise).

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- If Fourier's law does not hold,  $\kappa_N \sim N^\delta$ , universality of the diverging order  $\delta$  of the conductivity?
- Numerical simulations are not conclusive ( $\delta \in [0.25; 0.47]$  for the same models) and subject of intense debate.

- FPU chains are mathematically very difficult to study. We perturb the Hamiltonian dynamics by a stochastic noise. These stochastic perturbations simulate (qualitatively) the long time (chaotic) effect of the deterministic nonlinear model.



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Noise 2 = energy and momentum conservative

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# Construction of the noise

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- Noise 2,  $d \geq 1$  ... are of the same type.

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*Inconvenient* : The quantity to evaluate is a dynamical quantity.

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with  $M$  martingale (mean 0 w.r.t. any initial condition)

$$j_{x, x+e_k} = -\frac{1}{2}(\nabla V)(q_{x+e_k} - q_x) \cdot (p_{x+e_k} + p_x) - \gamma \nabla_{e_k} p_x^2$$

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$\kappa_N \sim \kappa_{GK}$  with truncation of the time up to time  $t_N = N$  (phononic picture).

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#### Theorem (B., Olla, JSP'05)

- $\kappa_{GK}$  is finite (pinned or unpinned) in any dimension.
- Fourier's law holds and linear response theory is correct: System of length  $N$  in contact with two Langevin baths at temperature  $T_\ell$  and  $T_r$  in its steady state

$$\lim_{N \rightarrow \infty} N \langle j_{x,x+1} \rangle_{ss} = \kappa_{GK} (T_r - T_\ell)$$

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### Theorem (B., '08)

*Harmonic system with random masses and energy conservative noise 1. The conductivity defined by Green-Kubo formula is strictly positive and bounded above:*

$$0 < c_- \leq \kappa_{KG} \leq C_+ < +\infty$$

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Theorem (Basile, B., Olla, PRL'06)

$$C_{1,1}(t) = \lim_{N \rightarrow \infty} \left\langle \left( \sum_x j_{x, x+e_1}(t) \right), j_{0, e_1}(0) \right\rangle_{eq.}$$

$$C_{1,1}(t) = \frac{T^2}{4\pi^2 d} \int_{[0,1]^d} (\partial_{k^1} \omega(k))^2 e^{-t\gamma\psi(k)} dk$$

where  $\omega$  is the dispersion relation of the harmonic chain

$$\omega(k) = (\nu + 4\alpha \sum_{j=1}^d \sin^2(\pi k^j))^{1/2}, \quad \psi(k) \sim k^2$$

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### Corollary

- $C_{1,1}(t) \sim t^{-d/2}$  in the unpinned case ( $\nu = 0$ )
- $C_{1,1}(t) \sim t^{-d/2-1}$  in the pinned case ( $\nu > 0$ )

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### Corollary

If the system is unpinned ( $\nu = 0$ ) then "truncated" Green-Kubo formula for  $\kappa_N$  gives:

$$\begin{cases} \kappa_N \sim N^{1/2} & \text{if } d = 1 \\ \kappa_N \sim \log N & \text{if } d = 2 \end{cases}$$

In all other cases  $\kappa_N$  is bounded in  $N$  and converges to  $\kappa_{GK}$ .

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## Theorem (Basile,B.,Olla'08)

- For  $d \geq 3$ , if  $W > 0$  is "general" or if  $W = 0$  and  $0 < c_- \leq V'' \leq C_+ < +\infty$  then

$$\kappa_N \leq C.$$

- For  $d = 2$ , if  $W = 0$  and  $0 < c_- \leq V'' \leq C_+ < \infty$

$$\kappa_N \leq C(\log N)^2.$$

- For  $d = 1$ , if  $W = 0$  and  $0 < c_- \leq V'' \leq C_+ < \infty$ , then

$$\kappa_N \leq C\sqrt{N}.$$

- In any dimension, if  $V$  are quadratic and  $W > 0$  is "general" then  $\kappa_N \leq C$ .

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- Simulations ( $d = 1$ ) for unpinned systems with energy/momentum conservative noise 2. The strength of the noise is regulate by  $\gamma$ . Then  $\kappa_N \sim N^\delta$ .

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Kinetic approach :  $\delta = 2/5$  (Perverzev, Lukkarinen and Spohn)

MCT approach :  $\delta = 1/2$  (Delfini, Lepri, Livi, Politi)

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- Function  $h_\lambda$  is local in the energy conservative case and non-local in the energy/momentum case. For  $d \geq 3$  or  $\nu > 0$ , the decay of  $h_\lambda$  is sufficient to assure a finite conductivity.

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- If  $\delta T = T_r - T_\ell$  is small

$$f_{ss} = \mathbf{1} + \delta T h_0 + o((\delta T)^2)$$

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$$(N^{-1} - \gamma S)^{-1} \left( \sum_x j_{x, x+e_1} \right) = \sum_{j=1}^d \sum_{x, y} G_N(x-y) p_x^j V'(q_{y+e_1}^j - q_y^j)$$

where  $G_N(z)$  is the solution of the resolvent equation

$$N^{-1} G_N(z) - 2\gamma(\Delta G_N)(z) = -\frac{1}{2} [\delta_0(z) + \delta_{e_1}(z)]$$