A contro problem

Partiallyobserved Langevin equation

Examples

Model reduction of partially-observed stochastic differential equations

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Outline

Motivation

A control problem

Partiallyobserved Langevin equation

Examples

1 Motivation

Conformational flexibility

2 A control problem

- Balanced truncation
- Port-controlled Hamiltonian systems

3 Partially-observed Langevin equation

- Controllability and observability
- Model reduction by balancing

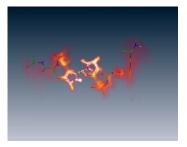
Motivation

A control problem

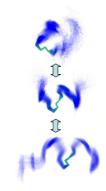
Partiallyobserved Langevin equation

Examples

Biological function depends on conformation's flexibility.



 $1.3\mu s$ simulation of dodeca-alanine at T = 300K (implicit solvent, GROMOS96 force field)



Conformational flexibility

 The microscopic dynamics (molecule & solvent) is generated by the nonlinear Hamiltonian

$$H: T^*Q \to \mathbb{R}, Q \subseteq \mathbb{R}^{3n}, H = p^T M^{-1}p + V(q),$$

with initial conditions distributed according to exp(−βH).
We suppose that the dynamics within a conformation can be approximated by the linear Langevin equation

$$\begin{split} \dot{x}(t) &= (J - D) \, \nabla \mathcal{H}_{\text{lin}}(x(t)) + S \dot{\mathcal{W}}(t) \\ \text{where } \mathcal{H}_{\text{lin}} &= \frac{1}{2} x_2^T \bar{\mathcal{M}}^{-1} x_2 + \frac{1}{2} x_1^T \bar{\mathcal{L}} x_1 \text{ and} \\ J &= \begin{pmatrix} \mathbf{0} & \mathbf{1} \\ -\mathbf{1} & \mathbf{0} \end{pmatrix}, \, D = \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \gamma \end{pmatrix}, \, S = \begin{pmatrix} \mathbf{0} \\ \sigma \end{pmatrix} \end{split}$$

Motivation

A contro problem

Partiallyobserved Langevin equation

Examples

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Conformational flexibility, cont'd

Motivation

A contro problem

Partiallyobserved Langevin equation

- For γ s.p.d, the system is **stable**, i.e., all eigenvalues of $A = (J D)\nabla^2 H_{\text{lin}}$ have strictly negative real part.
- The system satisfies **Hörmander's condition**. If further $2\gamma = \beta \sigma \sigma^{T}$, this entails ergodicity with respect to

$$d\mu(x) \propto \exp(-eta \mathcal{H}_{ ext{lin}}(x)) dx$$
 .

- The Gaussian distribution exp(-βH_{lin}) indicates that **all** modes are flexible. Which are the most flexible ones?
- Often the most flexible modes are thought of as having the largest variance. But: variance is not always most important to the dynamics.

A control problem

Partiallyobserved Langevin equation

Examples

1 Motivation

Conformational flexibility

2 A control problem

- Balanced truncation
- Port-controlled Hamiltonian systems

3 Partially-observed Langevin equation
 Controllability and observability
 Model reduction by balancing

Conformational flexibility as a control problem

Key observations

- Not all modes are equally stiff; moreover noise and friction may not be spatially isotropic.
- 2 Not all modes are observed (e.g., generalized momenta).
 - We may consider flexibility as the property of being sensitive to excitations due to the noise (controllability).
 - A sensible notion of flexibility should take into account what can be measured experimentally (observability).
 - Determining the flexibility of a conformation therefore amounts to identifying a low-dimensional subspace of easily controllable and highly observable modes

Motivation

A control problem

Partiallyobserved Langevin equation

Linear control systems

Motivation

A control problem

Partiallyobserved Langevin equation

Examples

■ For the moment let us replace the white noise by a smooth control function u ∈ L²([0, T]), i.e.,

$$\dot{x}(t) = Ax(t) + Su(t), \quad x(0) = 0$$

 $y(t) = Cx(t),$

where $A = (J - D)\nabla^2 H_{\text{lin}} \in \mathbb{R}^{2d \times 2d}$ and $y \in \mathbb{R}^k$ is a linear observable (e.g., all configurations $y = x_1$).

VoP yields the transfer function (input-output relation),

$$G: L^2([0,T]) \rightarrow \mathbb{R}^k, y(t) = C \int_0^t e^{A(t-s)} Su(s) ds.$$

Model reduction by balanced truncation

For the stable linear system $\dot{x} = Ax + Su$, y = Cx compute **controllability and observability Gramians**

$$Q = \int_0^\infty \exp(At)SS^T \exp(A^T t)dt$$
$$P = \int_0^\infty \exp(A^T t)C^T C \exp(At)dt$$

Balancing: find a transformation $x \mapsto Tx$, such that

$$T^{-1}QT^{-T} = T^{T}PT = \operatorname{diag}(\sigma_1, \ldots, \sigma_{2d}).$$

Truncation: Project onto the first *m* columns of *T*.

Motivation

A control problem

Partiallyobserved Langevin equation

Properties of balanced truncation

Interpretation of the controllability Gramian Q: x ∈ ℝ^{2d} is "more controllable" than x' ∈ ℝ^{2d} if

$$x^T Q x > x'^T Q x'$$
 ($|x| = |x'| = 1$).

Interpretation of the observability Gramian P: given an initial state x(0) = x and zero input, u = 0, we have

$$\|y\|_{L^2}^2 = x^T P x.$$

• **Approximation error** (H^{∞} error bound):

$$\sigma_{m+1} < \max_{u} \frac{\|(G - G_{trc})u\|_{L^2}}{\|u\|_{L^2}} < 2(\sigma_{m+1} + \ldots + \sigma_{2d}).$$

Glover 1984, Rowley 2005

Motivation

A control problem

Partiallyobserved Langevin equation

Port-controlled Hamiltonian systems

Let's go back to our second-order Langevin problem and consider the stable system

$$\begin{aligned} \dot{x}(t) &= (J-D)\nabla H_{\text{lin}}(x(t)) + Su(t) \\ y(t) &= Cx(t). \end{aligned}$$

- Balancing mixes configurations and momenta. Truncation (e.g., by projection) does not preserve structure.
- Preserving structure requires to impose constraints on the Hamiltonian part (energy & structure matrix). Then

$$\dot{\xi}(t) = (J_{\mathrm{trc}} - D_{\mathrm{trc}}) \nabla H_{\mathrm{trc}}(\xi(t)) + S_{\mathrm{trc}}u(t)$$

 $y(t) = C_{\mathrm{trc}}\xi(t)$

is stable with $\xi \in \mathbb{R}^m$ (odd or even dim.) and $J_{\text{trc}} = -J_{\text{trc}}^T$.

Hartmann et al. 2007

Motivation

A control problem

Partiallyobserved Langevin equation

A contro problem

Partiallyobserved Langevin equation

Examples

1 Motivation

Conformational flexibility

- 2 A control problem
 - Balanced truncation
 - Port-controlled Hamiltonian systems

3 Partially-observed Langevin equation

- Controllability and observability
- Model reduction by balancing

Partially-observed Langevin equation

• Consider the family of Langevin equations for $\epsilon > 0$

$$\begin{aligned} \dot{x}^{\epsilon}(t) &= (J-D) \nabla \mathcal{H}_{\text{lin}}(x^{\epsilon}(t)) + \sqrt{\epsilon} S \dot{\mathcal{W}}(t) \\ y^{\epsilon}(t) &= C x^{\epsilon}(t) \,. \end{aligned}$$

• Using again the shorthand $A = (J - D)\nabla^2 H_{\text{lin}}$, we have $Y_t^{\epsilon} = CX_t^{\epsilon}$, where X_t^{ϵ} , $X_0^{\epsilon} = 0$ is the family of solutions

$$X_t^{\epsilon} = \sqrt{\epsilon} \int_0^t e^{A(t-s)} S \, dW(s) \, .$$

The system is **stable for all** $\epsilon > 0$. If $2D = SS^T$ it admits an ergodic invariant measure $d\mu^{\epsilon} \propto \exp(-H_{\text{lin}}/\epsilon)$.

Motivatior

A contro problem

Partiallyobserved Langevin equation

Controllability of the Langevin equation

Consider the map $F : H^1([0, T]) \to C([0, T])$ and f = F(u), u(0) = 0 that is defined point-wise by

$$(Fu)(t) = \int_0^t e^{A(t-s)} Su(s) ds.$$

We introduce the controllability function (rate function)

$$I_{x}(f) = \inf_{u \in H^{1}, f(T)=x} \int_{0}^{T} |\dot{u}(t)|^{2} dt$$

measuring the **minimum "energy"** that is needed to steer the system from f(0) = 0 to f(T) = x within time T.

• We declare $I_x(f) = \infty$ if no such $u \in H^1$ exists.

cf. Dembo & Zeitsouni 1998

Motivation

A contro problem

Partiallyobserved Langevin equation

Controllability of the Langevin equation, cont'd

Theorem (Hartmann, Schütte 2008)

The controllability function $L_{con}(x) = I_x(f)$ is given by

$$L_{\rm con}(x) = x^T Q(T)^{-1} x$$

with $Q(T) = cov(X^{\epsilon}(T))$ for $\epsilon = 1$.

- Proof: Minimize $||u||_{H^1}^2$ subject to (Fu)(T) = x.
- The idea of replacing the Brownian motion W(t) by its polygonal approximation $u \in H^1$ is to make sense of $I_x(f)$. If ϵ is small, LDT guarantees that f(t) is "close" to $X^{\epsilon}(t)$.
- For T → ∞, the controllability Gramian Q is the unique s.p.d solution of the Lyapunov equation

$$AQ + QA^T = -SS^T$$
.

Motivation

A contro problem

Partiallyobserved Langevin equation

Observability of the Langevin equation

The observability function

 $L_{\rm obs}(x) = \int_0^T |Y_t^0|^2 dt \,, \quad X_0^0 = x$

measures the **output energy** up to time T in the absence of noise (i.e., $\epsilon = 0$), if $X_0^0 = x$ at time t = 0.

• For $T \to \infty$, it follows immediately that

$$L_{\rm obs}(x) = x^T P x$$

where the **observability Gramian** P is the unique s.p.d solution of the Lyapunov equation

$$A^T P + P A = -C^T C.$$

Hartmann & Schütte 2008, cf. Moore 1981

Motivation

A contro problem

Partiallyobserved Langevin equation

Model reduction by balancing

- Compute the Gramians Q, P of the stable Langevin system $\dot{x} = (J D)\nabla H(x) + SW$, y = Cx.
 - Find the balancing transformation $x \mapsto Tx$, such that

$$T^{-1}QT^{-T} = T^{T}PT = \operatorname{diag}(\sigma_1, \ldots, \sigma_{2d}).$$

Notice:
$$T^{-1}QPT = \text{diag}(\sigma_1^2, \ldots, \sigma_{2d}^2).$$

Constrain the system to the subspace of the largest singular values σ₁,..., σ_m or, alternatively, scale the smallest Hankel singular values according to

$$(\sigma_{m+1},\ldots,\sigma_{2d})\mapsto \delta(\sigma_{m+1},\ldots,\sigma_{2d}), \quad \delta>0.$$

and balance the Langevin equation by $x \mapsto T_{\delta}x$. In the resulting perturbed system, let δ go to zero.

Motivation

A contro problem

Partiallyobserved Langevin equation

Properties of the balanced Langevin equation

 Restriction to the best controllable and observable subspace yields the constrained Langevin equation

$$\dot{\xi}(t) = (J_{
m trc} - D_{
m trc}) \nabla H_{
m trc}(\xi(t)) + S_{
m trc} \dot{W}(t)$$

 $y(t) = C_{
m trc} \xi(t)$

with $J_{\rm trc} = -J_{\rm trc}^{T}$ and $H_{\rm trc}$ as before.

■ By 2D_{trc} = βS_{trc}S^T_{trc} and **asymptotic stability** it follows that the dynamics has the ergodic invariant measure

$$d
ho \propto \exp(-eta H_{
m trc})$$
 .

• Convergence of the scaled system for $\delta \rightarrow 0$ is due to **singular perturbation arguments**. In this case the effective Hamiltonian equals the free energy

$$H_{\rm free} = -\beta^{-1} \ln \rho(\xi) \,.$$

Hartmann et al. 2007, cf. Berglund & Gentz 2006, Kifer 2001

Motivation

A contro problem

Partiallyobserved Langevin equation

A contro problem

Partiallyobserved Langevin equation

Examples

1 Motivation

Conformational flexibility

- 2 A control problem
 - Balanced truncation
 - Port-controlled Hamiltonian systems



- Partially-observed Langevin equation
 Controllability and observability
 - Model reduction by balancing

Unobservable modes

Consider the second-order high-friction Langevin equation

$$\begin{aligned} \ddot{x}_1 &= -\epsilon x_1 - \dot{x}_1 + \sqrt{2\epsilon} \dot{\mathcal{W}} \\ y &= x_1 \,. \end{aligned}$$

Computing controllability and observability Gramian

$$Q = \left(egin{array}{cc} 1 & 0 \\ 0 & \epsilon \end{array}
ight), \quad P = rac{1}{2\epsilon} \left(egin{array}{cc} 1+\epsilon & 1 \\ 1 & 1 \end{array}
ight)$$

yields the joint HSV $\sigma_1 \sim 1/\sqrt{\epsilon}$ and $\sigma_2 \sim \sqrt{\epsilon}$.

To lowest order in ε, the constrained equation for ξ = x₁ turns out to be the first-order diffusion

$$\dot{x}_1 = -\epsilon x_1 + \sqrt{2\epsilon} \dot{W} \,.$$

Motivation

A contro problem

Partiallyobserved Langevin equation

Influence of unobservable modes

Motivation

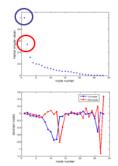
A contro problem

Partiallyobserved Langevin equation

Examples

Flexibility may come from unobservable modes: here, central dihedral angles and unobserved angular momenta of the freely rotating end-group are most important.

Helical conformation of octa-alanine, 14 dihedral angles plus conjugate momenta



Parametrization by HMMSDE: see Horenko & Hartmann 2007, Horenko & Schütte 2008

Take-home messages

Motivation

A contro problem

Partiallyobserved Langevin equation

- Mode balancing is a sensible approach towards molecular flexibility and model reduction that takes into account which variables can be observed.
- Efficient numerical tools for the computation of (exact or empirical) Gramians have recently become available.
- Structure-preservation for the reduced model is a subtle issue (e.g., singular structure matrix).
- Error bounds à la standard balanced truncation are still missing. Using averaging techniques may be a good idea.

A contro problem

Partiallyobserved Langevin equation

Examples

Thanks for your attention.

further information on biocomputing.mi.fu-berlin.de