

Fast irreversible Markov chains beyond the Metropolis paradigm

Werner Krauth

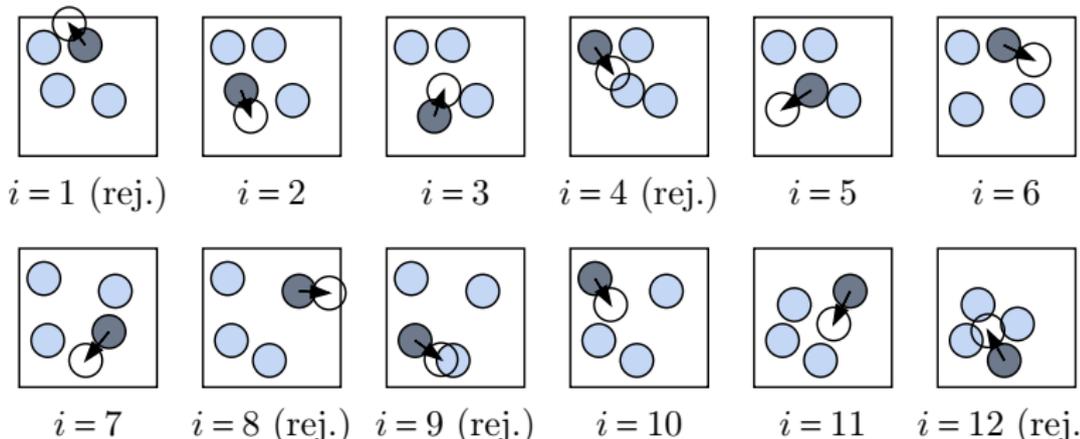
Laboratoire de Physique statistique
Département de Physique
Ecole normale supérieure
Paris, France

04 February 2016, Paris
COSMOS workshop

COmputational S tatistics and MOlecular S imulation

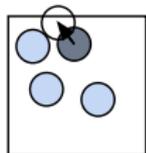
Collaborators: M. Michel, S. C. Kapfer,
M. Isobe, N. Yoshihiko, K. Hukushima,
E. P. Bernard, D. B. Wilson

- Local hard-sphere Monte Carlo:

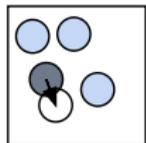


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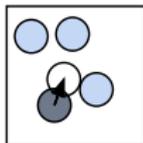
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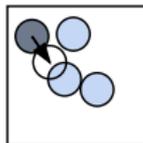
$i = 1$ (rej.)



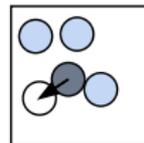
$i = 2$



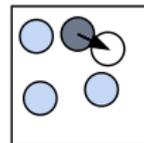
$i = 3$



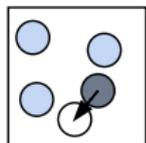
$i = 4$ (rej.)



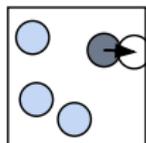
$i = 5$



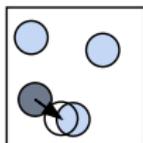
$i = 6$



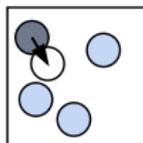
$i = 7$



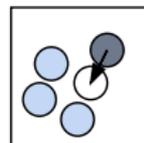
$i = 8$ (rej.)



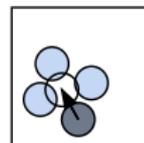
$i = 9$ (rej.)



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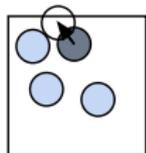
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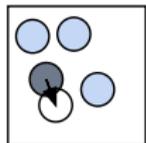
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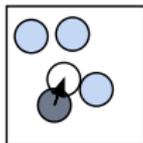
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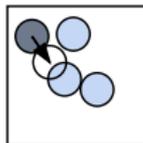
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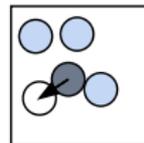
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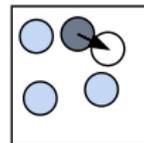
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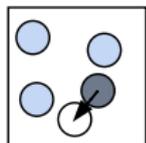
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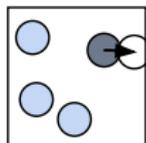
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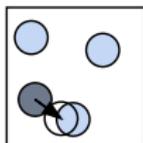
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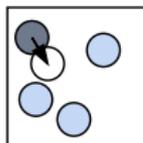
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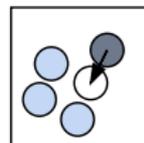
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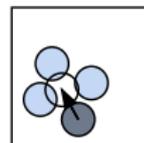
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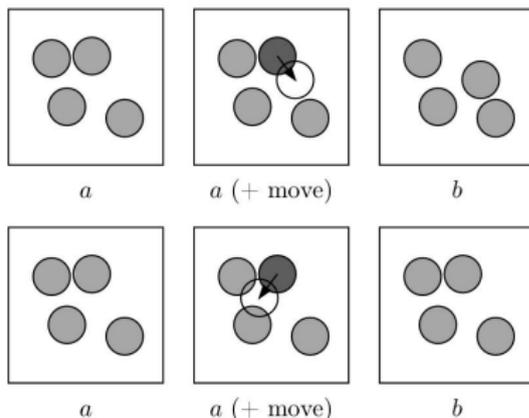
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Hard spheres

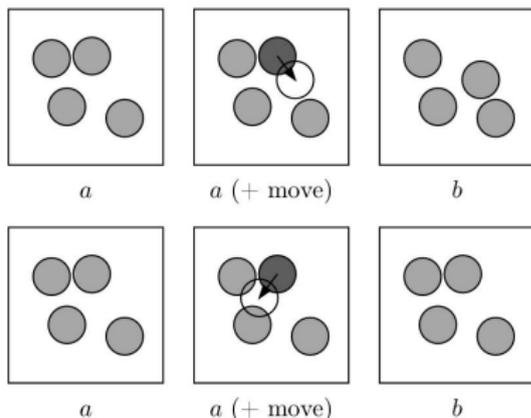


- Displacements (δ_x, δ_y) sampled uniformly in $([-\delta, \delta], [-\delta, \delta])$
- Algorithm satisfies **detailed balance**:

$$\pi(a)p(a \rightarrow b) = \pi(b)p(b \rightarrow a)$$

with $\pi(a) = \pi(b) = \pi(c)$, for all legal configurations of hard spheres

Hard spheres

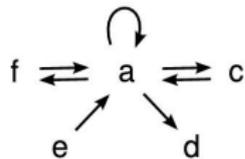


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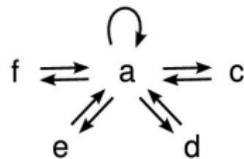
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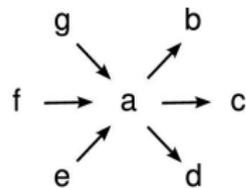
Detailed balance - global balance



global balance



detailed balance



maximal global balance

- flow in \equiv flow out (global balance condition):

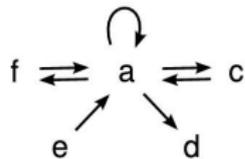
$$\underbrace{\sum_k \pi(k) p(k \rightarrow a)}_{\text{flow into } a} = \underbrace{\sum_k \pi(a) p(a \rightarrow k)}_{\text{flow out of } a}$$

- flow $\varphi(a \rightarrow b) \equiv$ flow $\varphi(b \rightarrow a)$ (detailed balance condition):
 - Metropolis algorithm (for flows and cond. probabilities)

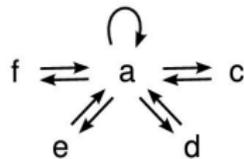
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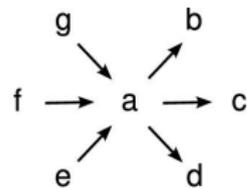
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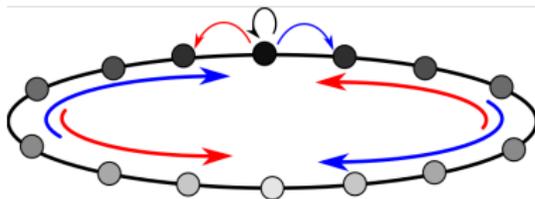
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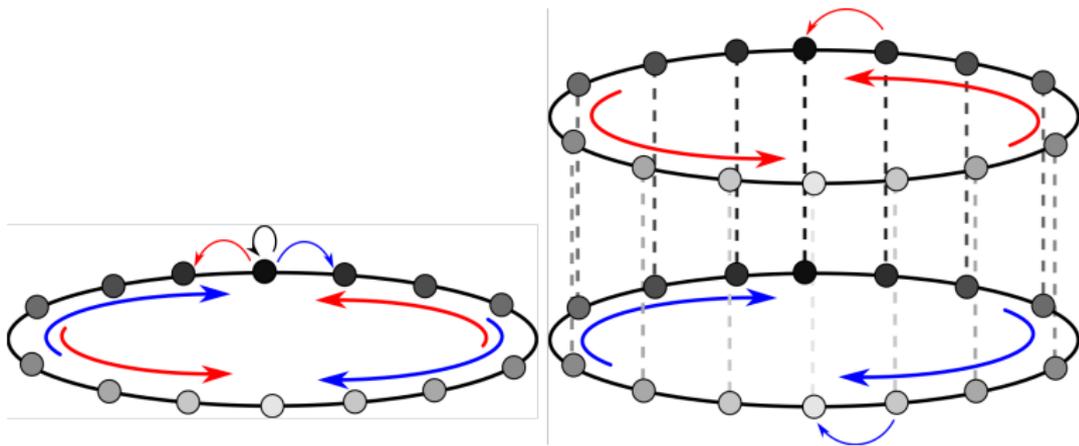
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Lifting - one (hard) sphere 1/2

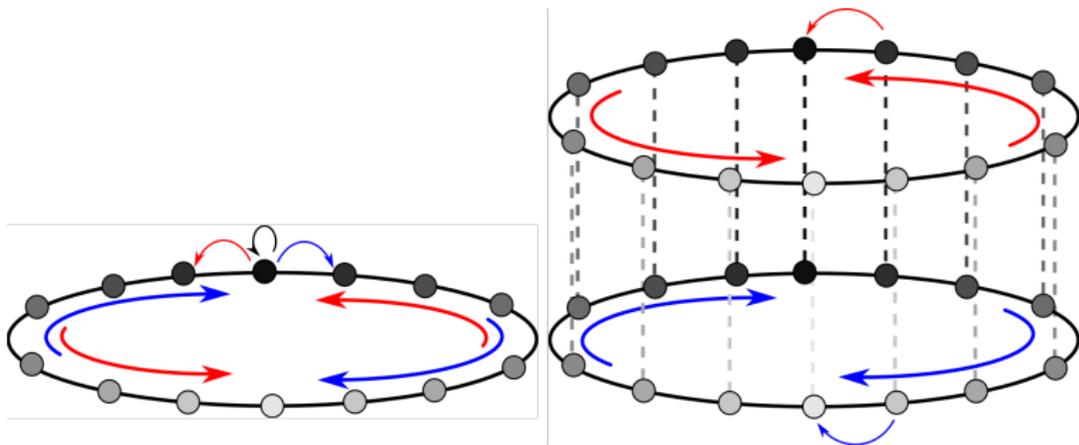


Lifting - one (hard) sphere 2/2



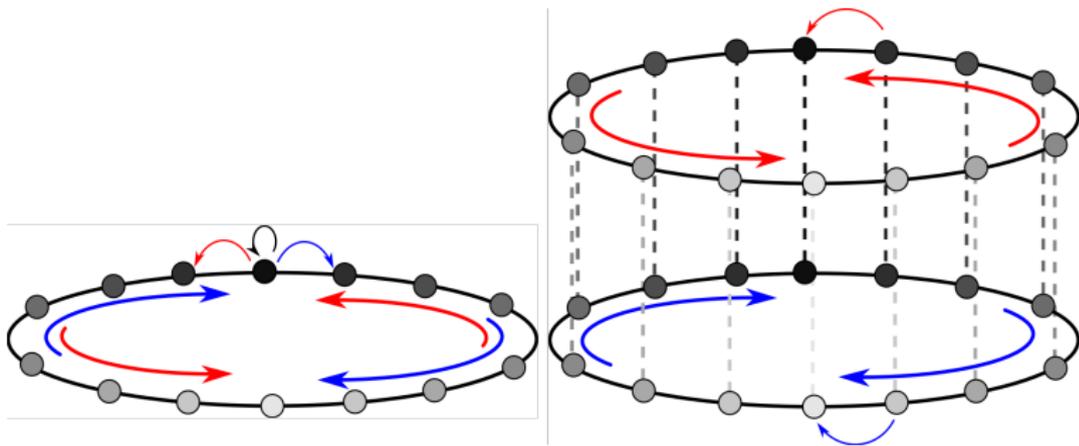
- Diaconis et al (2000)
- **lifting** \equiv additional variable
- Dynamical critical exponent $z = 1$ rather than $z = 2$
- Irreversible Markov chain

Lifting - one (hard) sphere 2/2



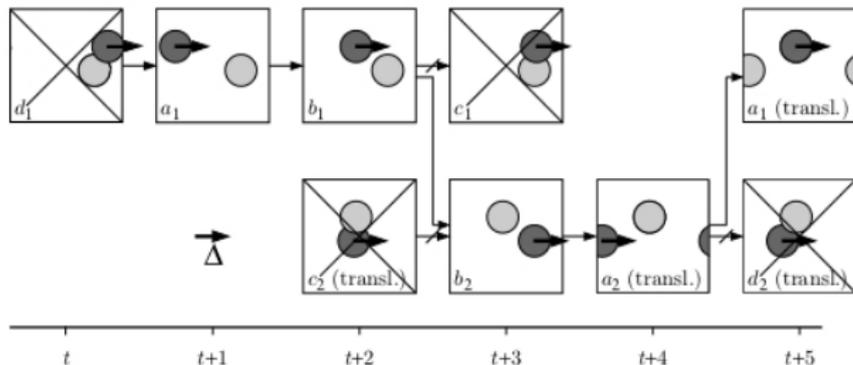
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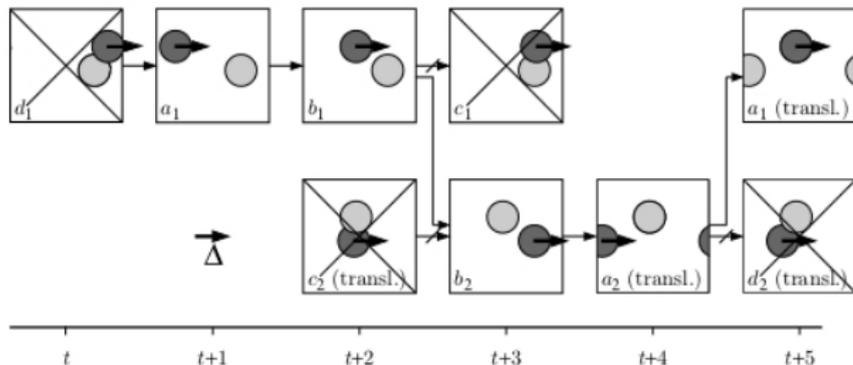
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Lifting - two hard spheres 1/1



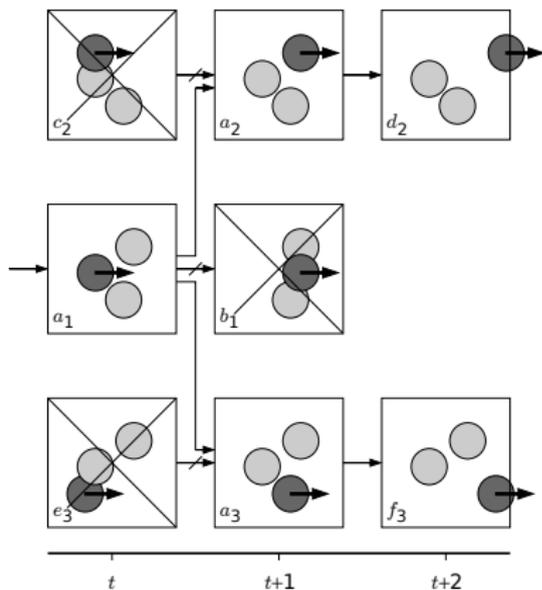
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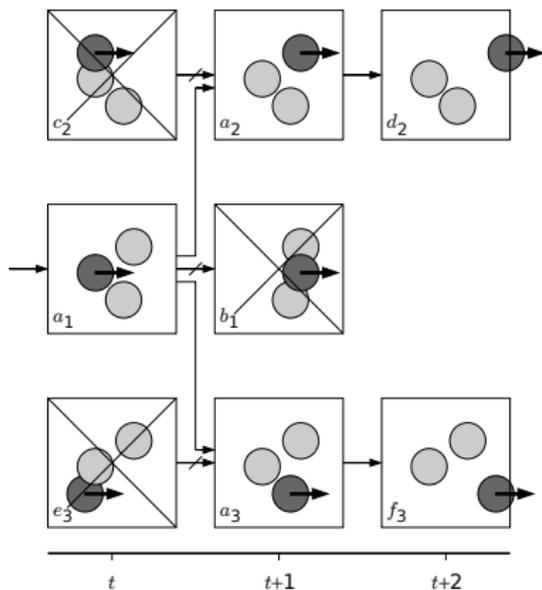
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Lifting - N hard spheres



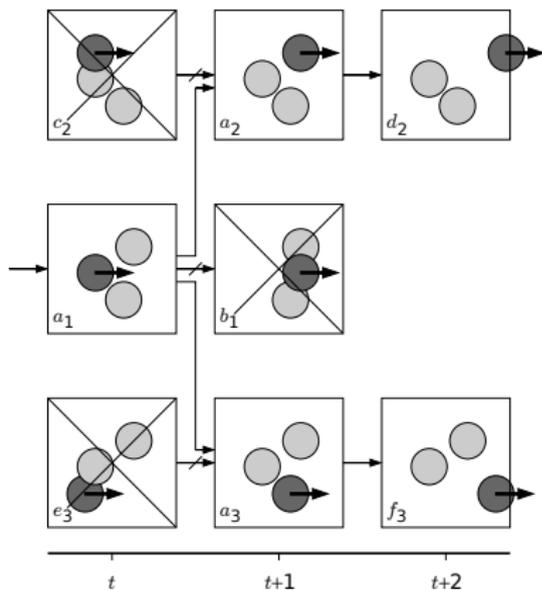
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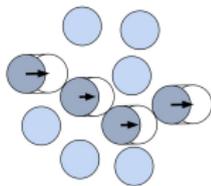
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Lifting algorithm for general potentials

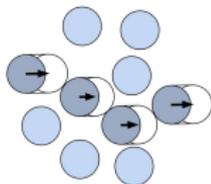


- Infinitesimal physical moves
- Liftings rather than rejections.
- Global balance rather than detailed balance.
- Bernard, Krauth, Wilson, (2009).

For general pair potentials, replace Metropolis by factorized filter

- $p^{\text{Met}}(a \rightarrow b) = \min \left[1, \exp(-\beta \sum_{i < j} (E_{ij}^b - E_{ij}^a)) \right]$
- $p^{\text{Met}}(a \rightarrow b) = \min \left[1, \prod_{i < j} \exp(-\beta (E_{ij}^b - E_{ij}^a)) \right]$
- $p^{\text{fact}}(a \rightarrow b) = \prod_{i < j} \min \left[1, \exp(-\beta (E_{ij}^b - E_{ij}^a)) \right]$
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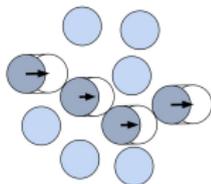


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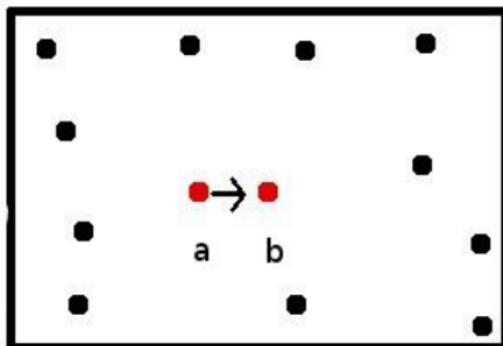


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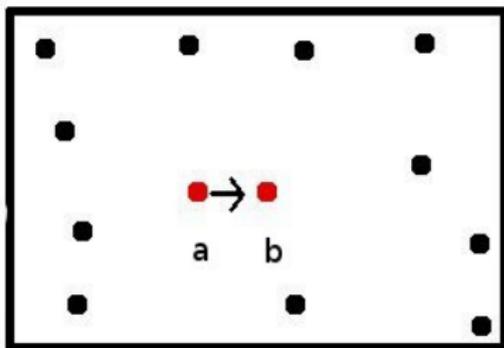
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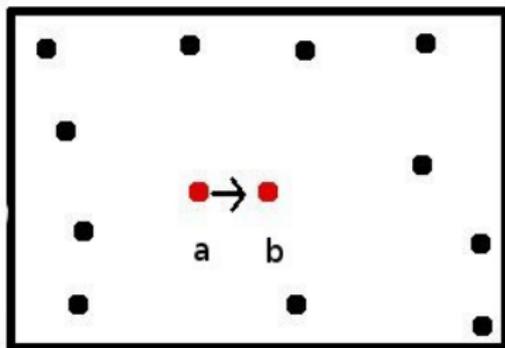
- Energy-based: MCMC knows its own weight.

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- pair-energy based
- bad idea, because $p^{\text{fact}}(a \rightarrow b) \leq p^{\text{Met}}(a \rightarrow b)$

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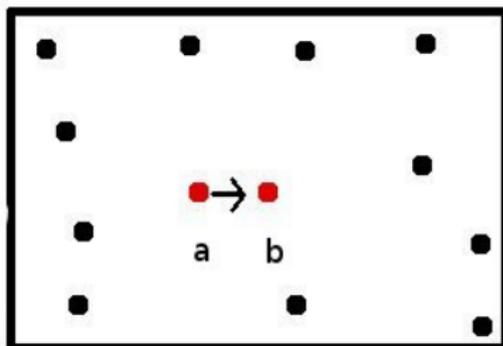


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Factorized filter 2/2

Acceptance probability

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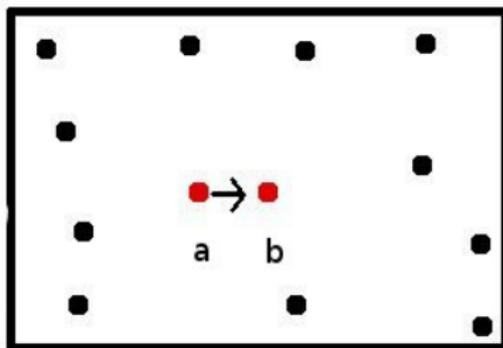


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- consensus rule (European-union like)
- Infinitesimal move: at most a single rejection
- lifting framework applicable

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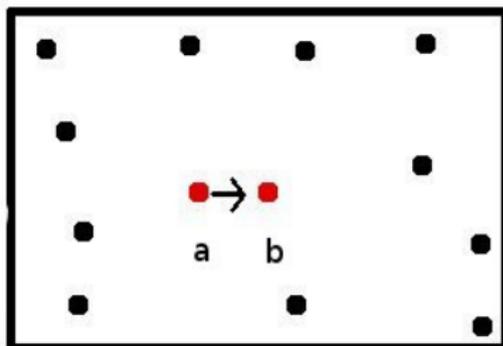


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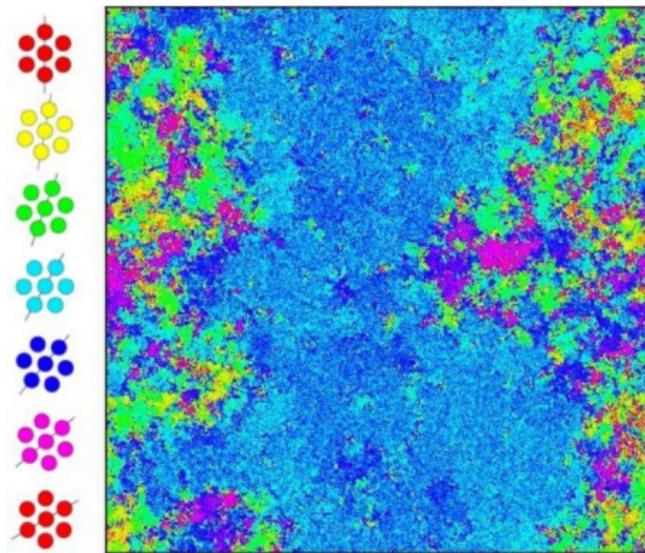
Factorized filter 2/2

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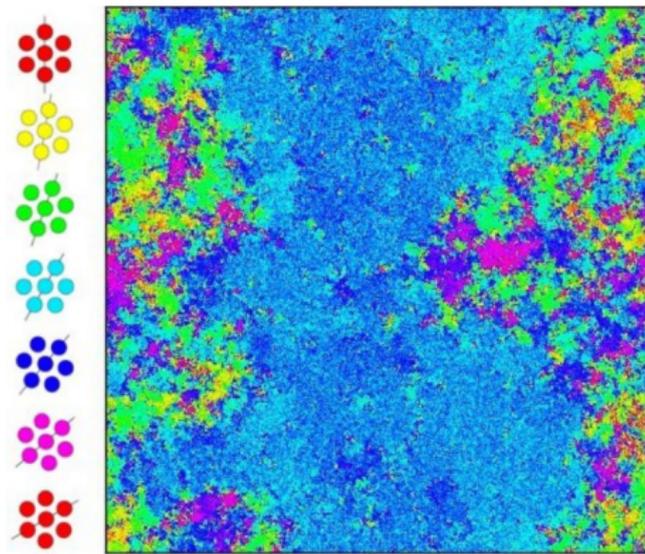
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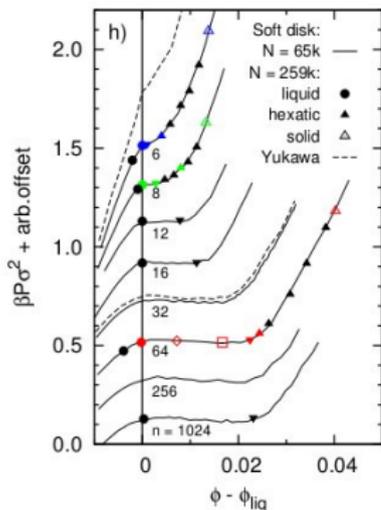
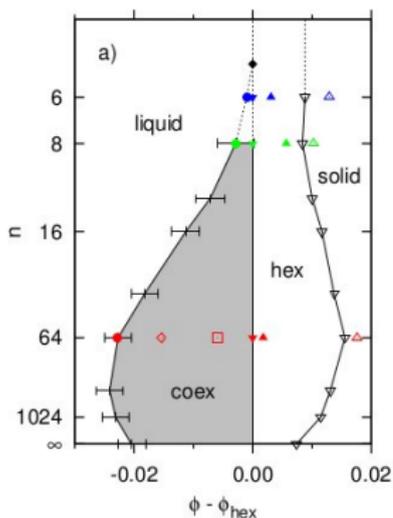
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- First-order liquid-hexatic transition in hard disks.



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Applications - Soft disks

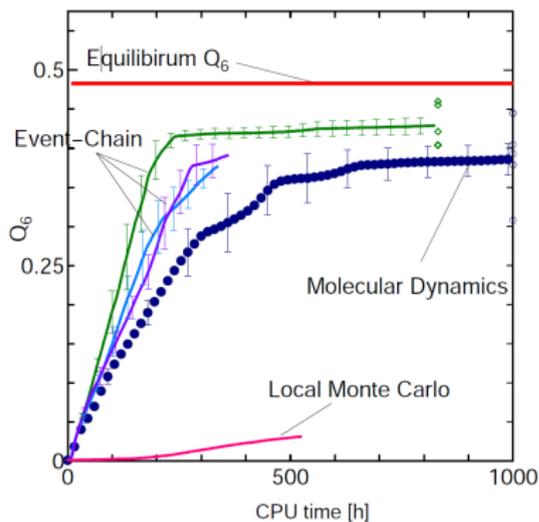
- Soft disks: $V \propto (\sigma/r)^n$.
- Mapping to Yukawa potentials.



- Kapfer & Krauth (PRL 2015).
- Two melting scenarios depending on softness n of potential.

Applications - Hard spheres

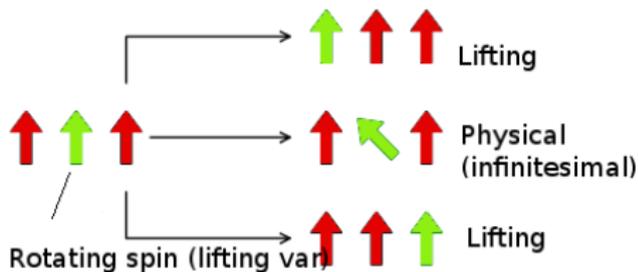
- Crystallization from liquid initial configurations ($\nu = 0.548$, $N = 1\text{Mio}$).



- Isobe & Krauth J. Chem. Phys. (2015)
- Considerable speedup of crystallization dynamics

Applications - Spin systems (XY, Heisenberg)

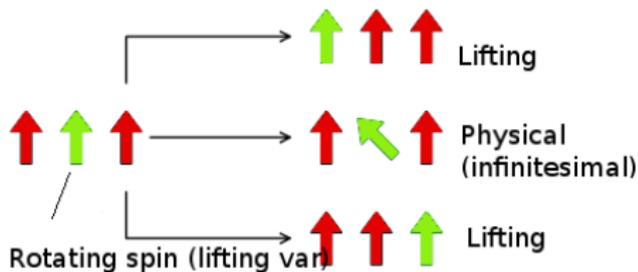
- Continuous spin systems:



- All rotations counter-clockwise, even for spin glass.
- Michel, Mayer & Krauth (EPL 2015)
fast, but $z \sim 2$ dynamical scaling in 2D XY model.
- Nishikawa, Michel, Krauth & Hukushima (PRE 2015),
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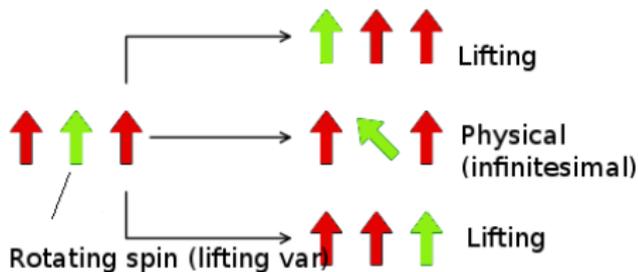
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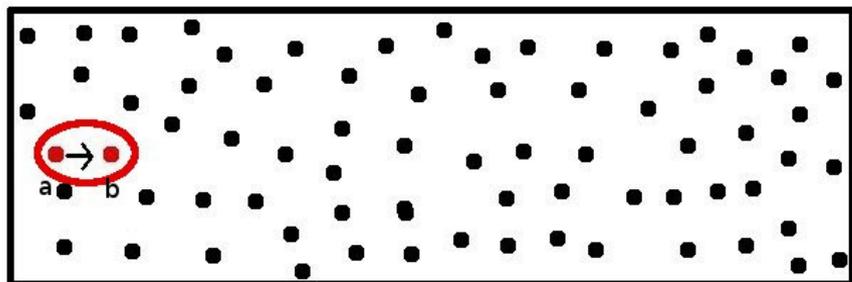
Applications - Spin systems (XY, Heisenberg)

- Continuous spin systems:



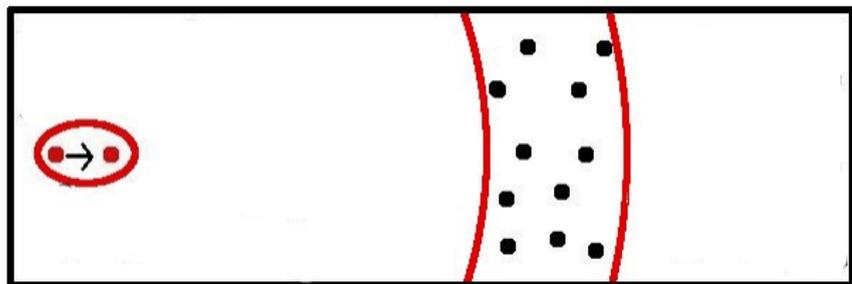
- All rotations counter-clockwise, even for spin glass.
- Michel, Mayer & Krauth (EPL 2015)
fast, but $z \sim 2$ dynamical scaling in 2D XY model.
- Nishikawa, Michel, Krauth & Hukushima (PRE 2015),
evidence for $z = 2 \rightarrow z = 1$ reduction in 3D Heisenberg model.

- $p^{\text{Met}}(a \rightarrow b) = \min [1, \exp(-\beta(E^b - E^a))]$
- Long-range particle system:



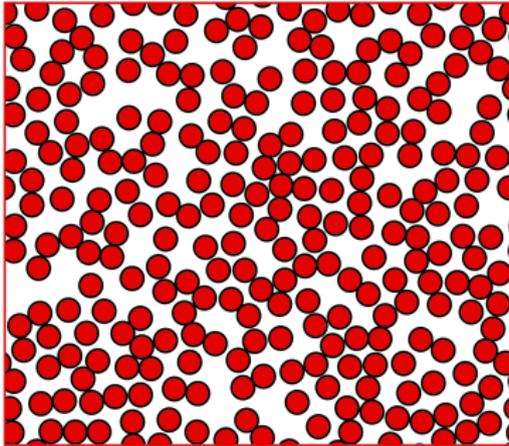
- Energy criterium problematic

- $p^{\text{fact}}(a \rightarrow b) = \prod_{i < j} \underbrace{\min \left[1, \exp(-\beta(E_{ij}^b - E_{ij}^a)) \right]}_{p_{ij}^{\text{accept}}}$

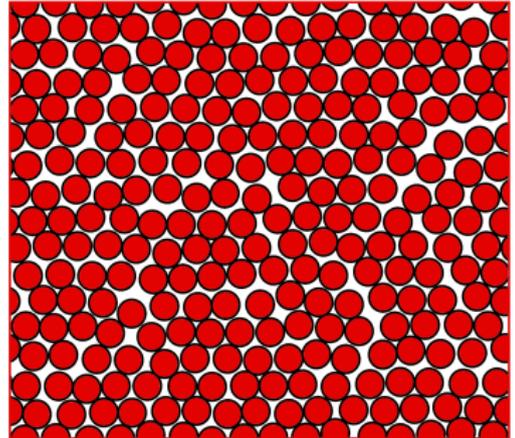


- if $p_{ij}^{\text{reject}} \ll 1$: consider subset of pairs.
- no more Ewald summation, no more Fourier methods.
- treat Coulomb forces directly.
- Kapfer & Krauth (manuscript in preparation).

2D melting transition



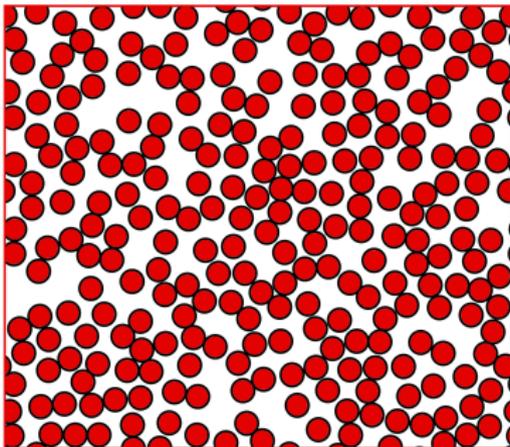
$$\eta = 0.48$$



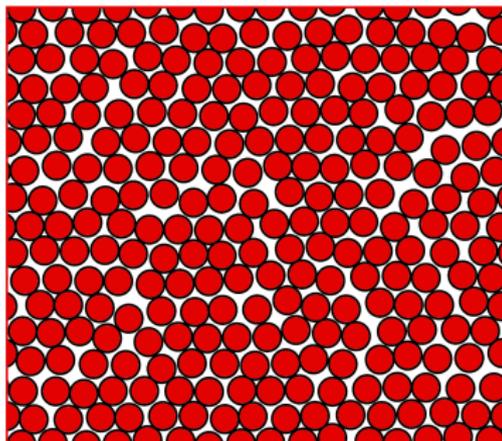
$$\eta = 0.72$$

- Generic 2D systems cannot crystallize (Peierls, Landau 1930s) but they can **turn solid** (Alder & Wainwright, 1962).
- Nature of transition disputed for decades (quid KTHNY?, quid hexatic?)

2D melting transition



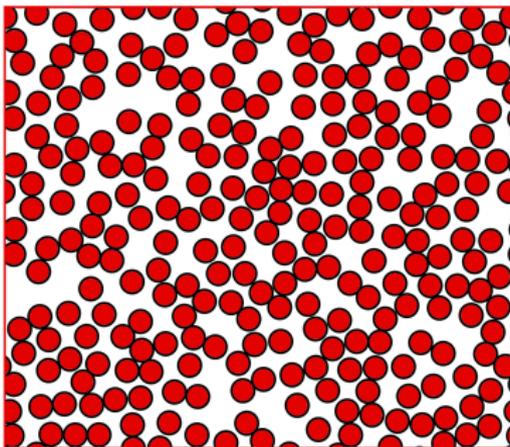
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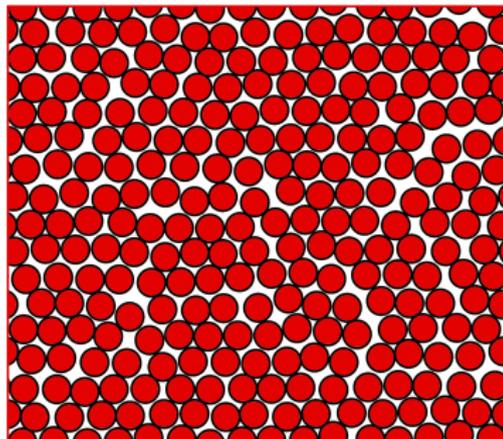
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Possible phases in two dimensions



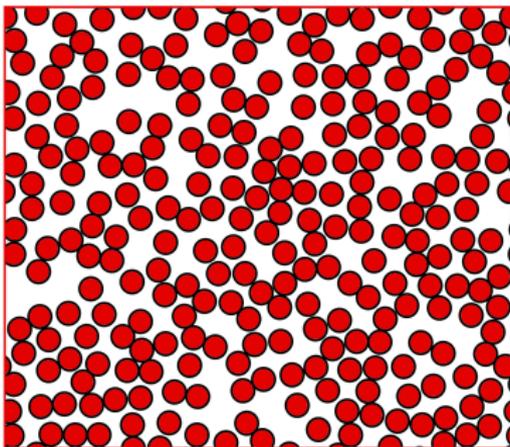
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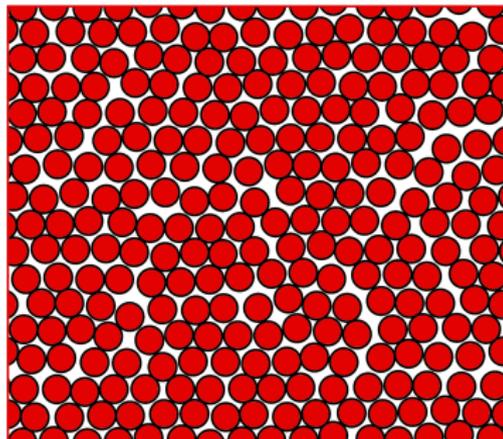
$$\eta = 0.72$$

Phase	positional order	orientational order
solid	algebraic	long-range
hexatic	short-range	algebraic
liquid	short-range	short-range

Possible phases in two dimensions



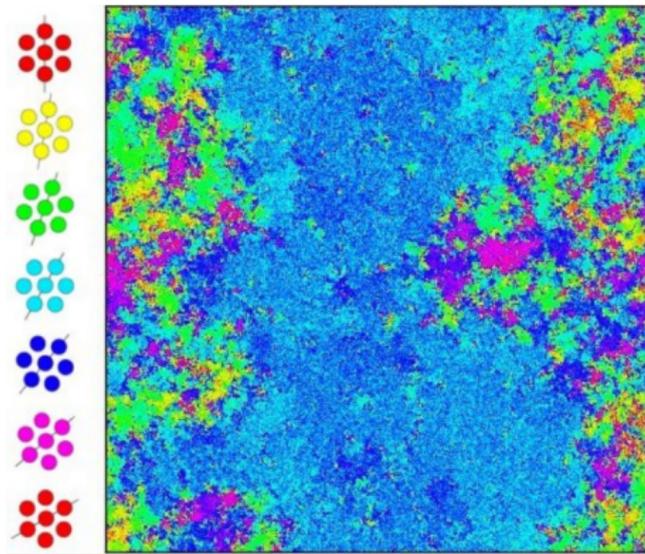
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$$\eta = 0.72$$

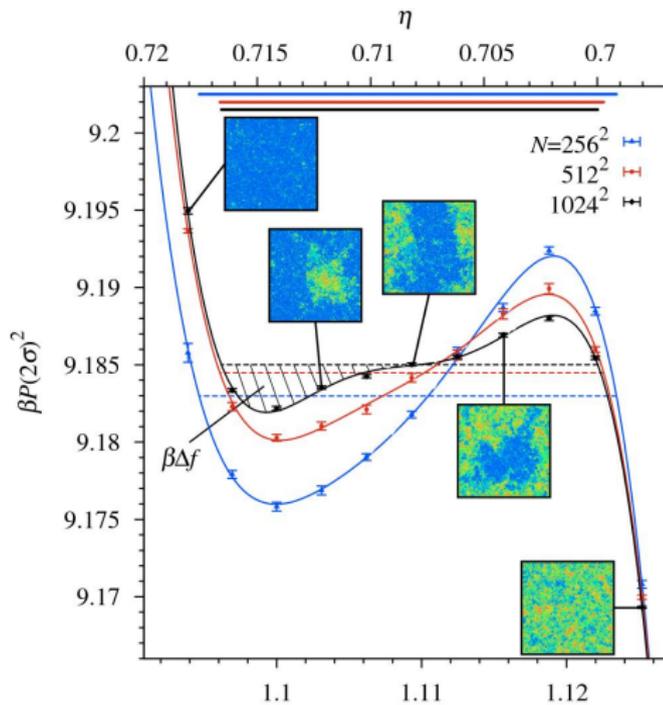
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Hard-disk configuration



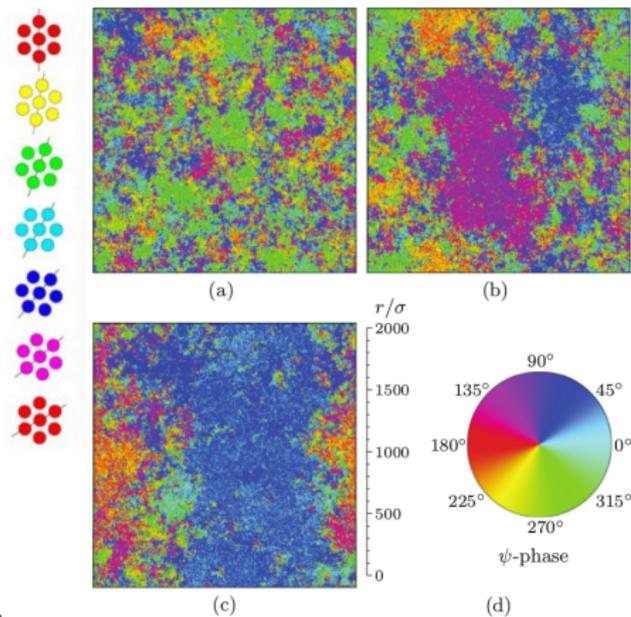
- 1024^2 hard disks
- circular color code for **orientational** order
- Bernard, Krauth (PRL 2011)

Equilibrium equation of state



- First-order transition (Bernard & Krauth, PRL (2011)).
- Many confirmations.

Phase coexistence in hard disks

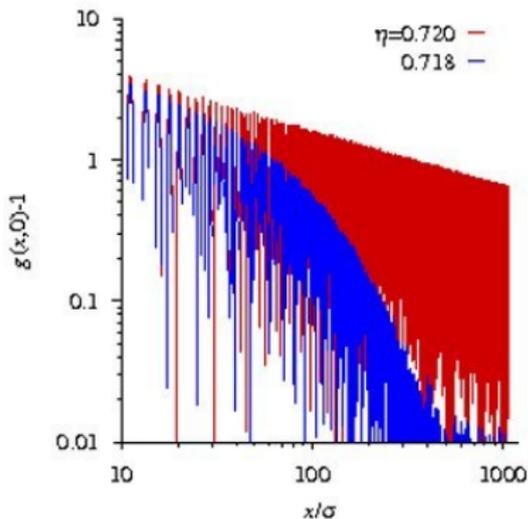
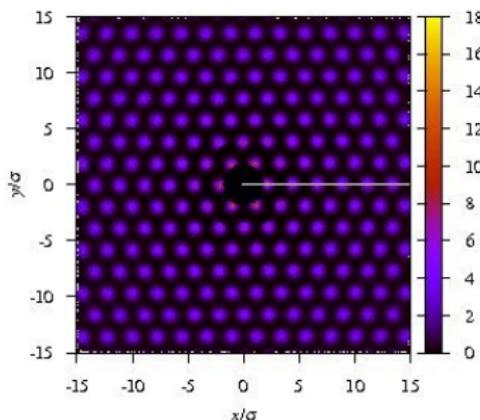


- 1024^2 systems.
- Densities $\eta = 0.700$ (a), $\eta = 0.704$ (b), $\eta = 0.708$ (c).
- Phase coexistence \implies Coarsening \implies Slow dynamics.
- cf. Engel et al (2013).

Possible phases (again)

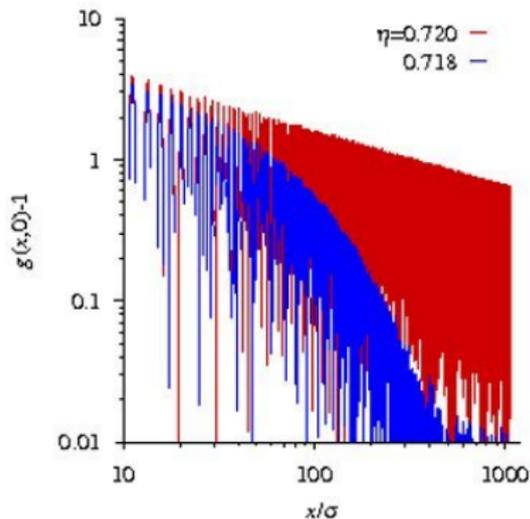
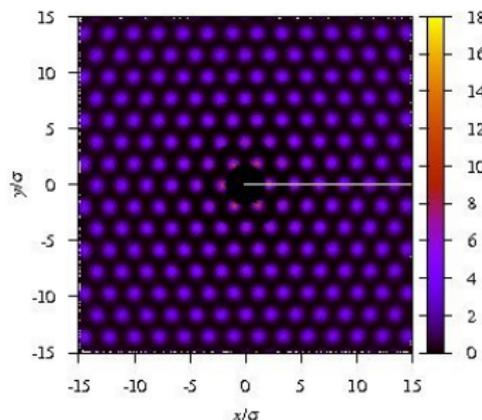
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Spatial correlations at $\eta = 0.718$ and 0.720



- Two-dimensional pair correlations, sample-averaged.
- At $\eta = 0.718$; hexatic.
- At $\eta \sim 0.720$: solid.
- Bernard & Krauth (PRL 2011).
- Many confirmations.

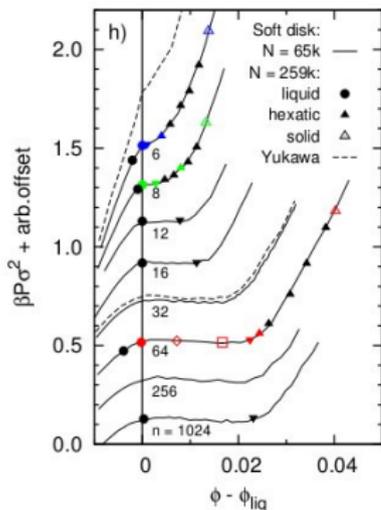
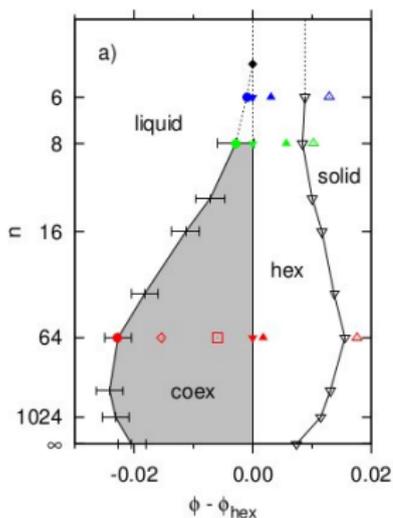
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- Many confirmations.

Applications - Soft disks

- Soft disks: $V \propto (\sigma/r)^n$.
- Mapping to Yukawa potentials.



- Kapfer & Krauth (PRL 2015).
- Two melting scenarios depending on softness n of potential.

- Event-chain, factorized Metropolis, lifting:
 - 'Beyond Metropolis' paradigm for Monte Carlo computations, mathematically challenging, completely general, many applications.
 - Makes infinitesimal moves.
 - Uses lifting.
 - Breaks detailed balance.
 - Ignores its own energy.
- Hard disks:
 - The mother of MCMC models & of 2D physics.
 - Hexatic phase exists, first-order liquid-hexatic transition.
 - Hexatic-solid transition is KT.
 - Communities A and B were wrong.
- Many extensions, both for physics and for algorithms.