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# Using Metropolis schemes to estimate correlation functions

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*Work also supported by ANR Funding ANR-14-CE23-0012 (“COSMOS”)*

Meeting GDR ISIS, Telecom (Paris), November 2015

# Motivation

- Computation of **integrated correlation functions**
  - transport coefficients in **molecular dynamics**
  - variance of time averages for SDEs  $\widehat{\varphi}_t = \frac{1}{t} \int_0^t \varphi(q_s) ds$
- Assume that...
  - the SDE  $q_t$  has a unique invariant measure  $\pi$
  - $\varphi$  has average 0 with respect to  $\pi$
  - time discretization with timestep  $\Delta t > 0 \rightarrow$  invariant measure  $\pi_{\Delta t}$

What is the numerical error arising from  $\Delta t > 0$ ?

$$\sigma_\varphi^2 = \lim_{t \rightarrow +\infty} t \mathbb{E}(\widehat{\varphi}_t^2) = 2 \int_0^{+\infty} \mathbb{E}(\varphi(q_t)\varphi(q_0)) dt$$

- Can be extended to the estimation of  $\int_0^{+\infty} \mathbb{E}(\varphi(x_t)\psi(x_0)) dt$

# Metropolize the discretization of the dynamics...?

- **Pros**

- Removes the **bias** on the invariant measure
- **Stabilizes** the discretization for non-globally Lipschitz drifts

- **Cons**

- Scaling of the rejection rate with the **dimension**
- Cannot be used for **non-reversible** dynamics...
- ... or worse: **nonequilibrium** systems for which the invariant measure is unknown!

- An early reference in the physics literature... “SmartMC” = MALA!

P. J. Rossky, J. D. Doll, and H. L. Friedman, Brownian dynamics as smart Monte Carlo simulation, *J. Chem. Phys.* (1978)

## Error estimates for MALA (1)

- Potential energy function  $V$ , **invariant measure**  $\nu(dq) = Z^{-1} e^{-\beta V(q)} dq$

Proposal move (recall  $\nabla V = (\partial_{q_1} V, \dots, \partial_{q_d} V)$ , dimension  $d$ )

$$\tilde{q}^{n+1} = \Phi_{\Delta t}(q^n, G^n) = q^n - \beta \Delta t \nabla V(q^n) + \sqrt{2\Delta t} G^n$$

- **Acceptance rate:** Metropolis-Hastings criterion

$$A_{\Delta t}(q^n, \tilde{q}^{n+1}) = \min \left( \frac{e^{-\beta V(\tilde{q}^{n+1})} T_{\Delta t}(\tilde{q}^{n+1}, q^n)}{e^{-\beta V(q^n)} T_{\Delta t}(q^n, \tilde{q}^{n+1})}, 1 \right),$$

where  $T_{\Delta t}(q, q') = \left( \frac{1}{4\pi\Delta t} \right)^{d/2} \exp \left( -\frac{|q' - q + \beta \Delta t \nabla V(q)|^2}{4\Delta t} \right)$

Markov chain encoded by a transition function

$$q^{n+1} = \Psi_{\Delta t}(q^n, G^n, U^n) = q^n + \mathbf{1}_{U^n \leq A_{\Delta t}(q^n, \Phi_{\Delta t}(q^n, G^n))} (\Phi_{\Delta t}(q^n, G^n) - q^n)$$

## Error estimates for MALA (2)

- Numerical scheme = **Markov chain** characterized by **transition operator**

$$P_{\Delta t}\varphi(q) = \mathbb{E}\left(\varphi(q^{n+1}) \mid q^n = q\right)$$

- Reference continuous dynamics  $dq_t = -\beta\nabla V(q_t) dt + \sqrt{2} dW_t$ 
  - leaves  $\nu$  invariant
  - generator  $\mathcal{L} = -\beta\nabla V(q)^T \nabla + \Delta$  (where  $\Delta = \partial_{q_1}^2 + \dots \partial_{q_N}^2$ )
  - recall that  $\frac{d}{dt}\mathbb{E}(\varphi(q_t)) = \mathbb{E}(\mathcal{L}\varphi(q_t))$

### $\Delta t$ -expansion of the evolution operator

$$P_{\Delta t}\varphi = \varphi + \Delta t \mathcal{A}_1\varphi + \Delta t^2 \mathcal{A}_2\varphi + \dots + \Delta t^{p+1} \mathcal{A}_{p+1}\varphi + \Delta t^{p+2} r_{\varphi, \Delta t}$$

- Weak order**  $p$  when  $\sup_{0 \leq n \leq T/\Delta t} \left| \mathbb{E}[\varphi(x^n)] - \mathbb{E}[\varphi(x_{n\Delta t})] \right| \leq C\Delta t^p$
- Satisfied if  $\mathcal{A}_k = \frac{\mathcal{L}^k}{k!}$  for all  $1 \leq k \leq p$

## Example: Euler-Maruyama, weak order 1 (dimension 1)

- Scheme  $q^{n+1} = \Phi_{\Delta t}(q^n, G^n) = q^n - \beta \Delta t V'(q^n) + \sqrt{2\Delta t} G^n$
- Note that  $P_{\Delta t}\varphi(q) = \mathbb{E}_G [\varphi(\Phi_{\Delta t}(q, G))]$

- Technical tool: **Taylor expansion**

$$\varphi(q + \delta) = \varphi(q) + \delta\varphi'(q) + \frac{1}{2}\delta^2\varphi''(q) + \frac{\delta^3}{6}\varphi^{(3)}(q) + \dots$$

- Replace  $\delta$  with  $\sqrt{2\Delta t} G - \beta\Delta t V'(q)$  and **gather in powers of  $\Delta t$**

$$\begin{aligned}\varphi(\Phi_{\Delta t}(q, G)) &= \varphi(q) + \sqrt{2\Delta t} G\varphi'(q) \\ &\quad + \Delta t \left( G^2\varphi''(q) - \beta V'(q)\varphi'(q) \right) + \dots\end{aligned}$$

- Taking **expectations w.r.t.  $G$**  leads to

$$P_{\Delta t}\varphi(q) = \varphi(q) + \underbrace{\Delta t \left( \varphi''(q) - \beta V'(q)\varphi'(q) \right)}_{=\mathcal{L}\varphi(q)} + O(\Delta t^2)$$

## Error estimates for MALA (3)

- For MALA, it can be shown that

$$P_{\Delta t}\varphi = \varphi + \Delta t \mathcal{L}\varphi + \Delta t^2 \mathcal{T}\varphi + \Delta t^{5/2} r_{\varphi, \Delta t}$$

(Fractional power of  $\Delta t$  is a signature of Metropolis...)

- An important ingredient is that the rejection rate is of order  $\Delta t^{3/2}$

$$\mathbb{E}_G \left| A_{\Delta t} \left( q, q - \beta \Delta t \nabla V(q) + \sqrt{2\Delta t} G \right) \right] - 1 + \Delta t^{3/2} \bar{\xi}(q) \right|^p \leq C_p \Delta t^{2p}$$

- For **compact** position spaces, geometric ergodicity can be proved

### Error estimates on integrated correlation functions

$$\int_0^{+\infty} \mathbb{E} \left( \varphi(q_t) \varphi(q_0) \right) dt = \Delta t \sum_{n=0}^{+\infty} \mathbb{E}_{\Delta t} \left( \varphi(q^n) \varphi(q^0) \right) + O(\Delta t)$$

**The error is determined by weak type expansions**

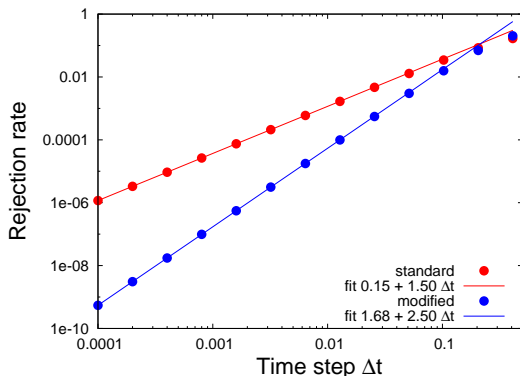
# Lower the rejection rate?

- Modifying the scheme to lower the rejection rate (1D expressions)

- modified drift  $-\beta V'(q) + \frac{\beta \Delta t}{6} (V^{(3)} - \beta V'' V')(q)$

- modified diffusion  $\text{Id} + \frac{\beta \Delta t}{3} V''(q)$

Rejection rate of order  $\Delta t^{5/2}$  but **weak order unchanged!**





## Modify the proposal functions

- Midpoint scheme: **implicit** hence more expensive...

$$\tilde{q}^{n+1} = q^n - \beta \Delta t \nabla V \left( \frac{\tilde{q}^{n+1} + q^n}{2} \right) + \sqrt{2\Delta t} G^n$$

- Hybrid Monte Carlo-like scheme

$$\tilde{q}^{n+1} = q^n - \beta \Delta t \nabla V \left( q^n + \frac{\sqrt{2\Delta t}}{2} G^n \right) + \sqrt{2\Delta t} G^n$$

Can be reformulated as (using  $h = \sqrt{2\beta\Delta t}$ )

$$p^n = \beta^{-1/2} G^n, \quad \begin{cases} q^{n+1/2} = q^n + \frac{h}{2} p^n, \\ p^{n+1} = p^n - h \nabla V \left( q^{n+1/2} \right), \\ \tilde{q}^{n+1} = q^{n+1/2} + \frac{h}{2} p^{n+1}. \end{cases}$$

**Reversible structure:** allows to compute the Metropolis ratio in terms of some extended energy difference  $H(q, p) = V(q) + p^2/2$

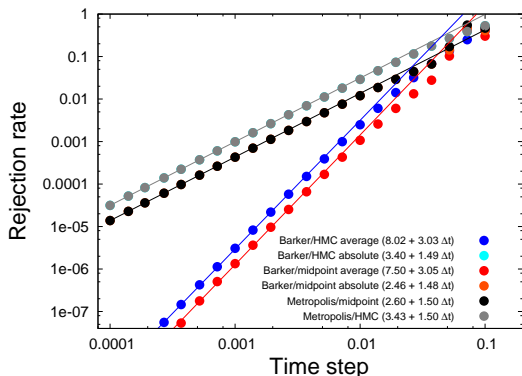
# Modify the acceptance criterion

- **Metropolis criterion**  $A_{\Delta t}^{\text{MH}}(q^n, \tilde{q}^{n+1}) = \min\left(1, e^{-\alpha \Delta t(q^n, \tilde{q}^{n+1})}\right)$

→ rejection rate  $1 - O(\Delta t^{3/2})$

- **Barker rule**  $A_{\Delta t}^{\text{Barker}}(q^n, \tilde{q}^{n+1}) = \frac{e^{-\alpha \Delta t(q^n, \tilde{q}^{n+1})}}{1 + e^{-\alpha \Delta t(q^n, \tilde{q}^{n+1})}}$

→ rejection rate  $1/2 + O(\Delta t^3)$  in average,  $1/2 + O(\Delta t^{3/2})$  in absolute value



# Results on integrated correlation functions

## Improved Green-Kubo formulas

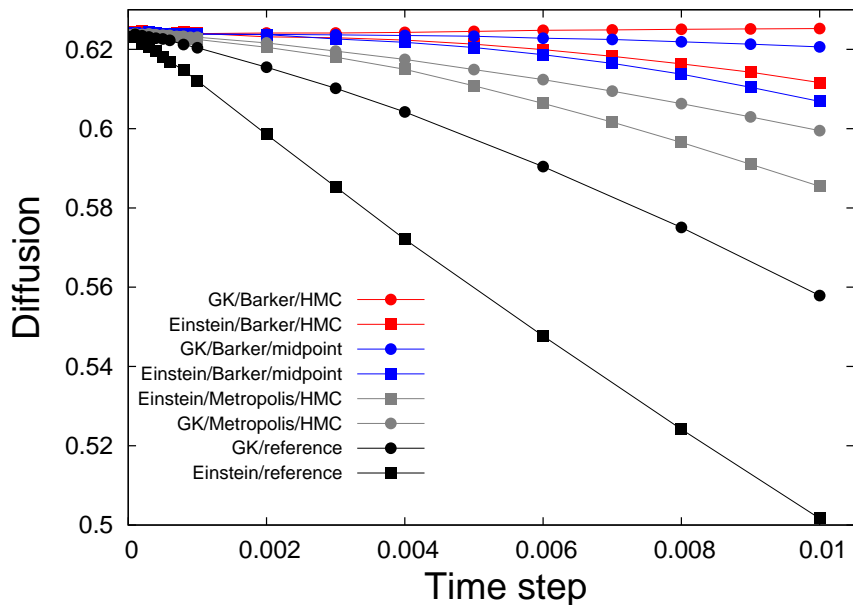
Set  $a = 1/2$  and  $\alpha = 2$  for Barker, and  $a = 1$  and  $\alpha = 3/2$  for Metropolis-Hastings. Then,

$$\int_0^{+\infty} \mathbb{E}[\psi(q_t)\varphi(q_0)] dt = \Delta t \left( a \sum_{n=0}^{+\infty} \mathbb{E}_{\Delta t}[\psi(q^n)\varphi(q^0)] - \frac{\mathbb{E}_\nu(\psi\varphi)}{2} \right) + O(\Delta t^\alpha)$$

### • Some comments...

- reduces to **trapezoidal** rule for Metropolis (but error  $\Delta t^{3/2}$ )
- **time renormalization** by a factor 2 for Barker
- statistical error increased by factor 2 for Barker, but reduced bias
- no fractional powers of  $\Delta t$  when Barker is used
- Key ingredient in the proof:  $\frac{P_{\Delta t} - \text{Id}}{\Delta t} \varphi = a \left( \mathcal{L}\varphi + \frac{\Delta t}{2} \mathcal{L}^2\varphi \right) + O(\Delta t^\alpha)$
- **Numerical illustration** for 1D system with  $V(q) = \cos(2\pi q)$  and  $\beta = 1$

# Results on integrated correlations $\varphi = \psi = V'$



# Conclusion and perspectives

- Numerical analysis of integrated correlation functions → bias
- Extension to dynamics with multiplicative noise

$$dq_t = \left( -\beta M(q_t) \nabla V(q_t) + \operatorname{div}(M)(q_t) \right) dt + \sqrt{2} M^{1/2}(q_t) dW_t$$

- Many open issues when the invariant measure is not known explicitly...  
→ **Nonequilibrium** systems in molecular dynamics

## References

- B. Leimkuhler, Ch. Matthews and G. Stoltz, The computation of averages from equilibrium and nonequilibrium Langevin molecular dynamics, *IMA J. Numer. Anal.* (2015)
- M. Fathi, A.-A. Homman and G. Stoltz, Error analysis of the transport properties of Metropolized schemes, *ESAIM Proc.* (2015)
- M. Fathi and G. Stoltz, Improving dynamical properties of stabilized discretizations of overdamped Langevin dynamics, *arXiv* **1505.04905** (2015)